

Narrow Escape Problems in 3D Domains

Daniel Gomez and Alexei F. Cheviakov

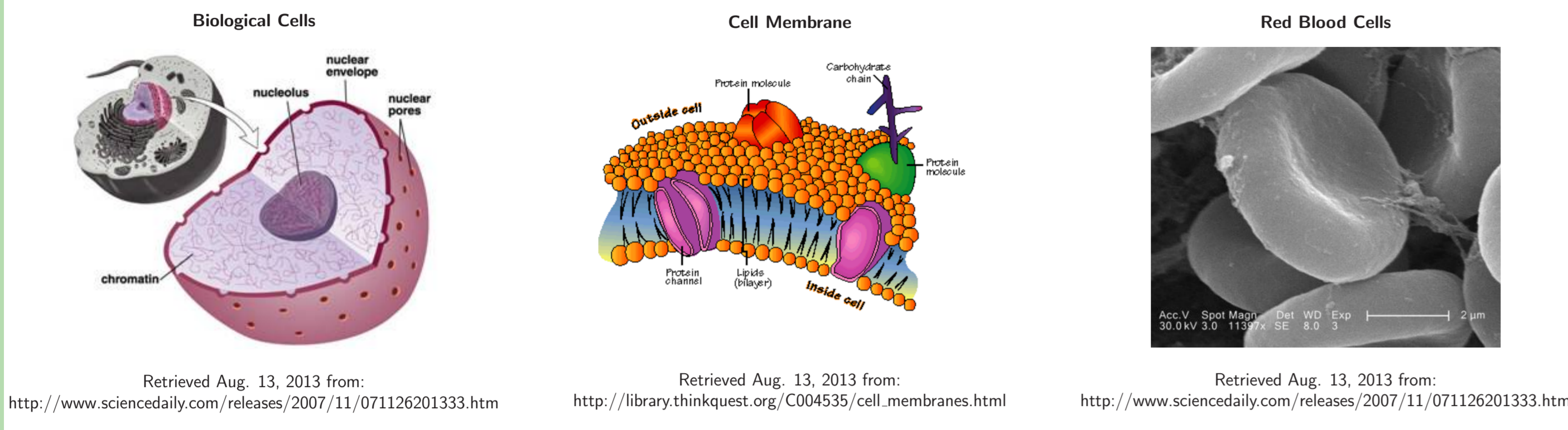
Department of Mathematics and Statistics, University of Saskatchewan

Motivation/Application

Numerous biological processes involve the transport of particles from a cell through its membrane:

- ▶ RNA transport through nuclear pores.
- ▶ Passive diffusion of molecules (e.g. CO₂ and O₂) through cell membrane.
- ▶ Diffusion of ions through protein channels (e.g. Na-K-Cl co-transporter in blood cells).

Typical size of transport regions is ~0.1% relative to overall cell size.



The Narrow Escape Problem

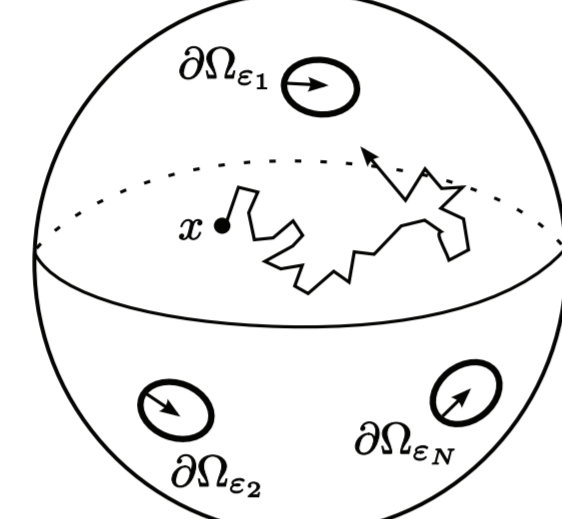
The Narrow Escape Problem (NEP) consists in finding the mean first passage time (MFPT) for a particle undergoing Brownian motion to escape an enclosing three-dimensional domain.

- ▶ Ω : three-Dimensional domain.
- ▶ $\partial\Omega_{\epsilon_j}$: absorbing boundary trap ($j = 1, \dots, N$).
- ▶ $v(\mathbf{x})$: MFPT for particle starting at $\mathbf{x} \in \Omega$.
- ▶ D : diffusion coefficient.

$$\bar{v} \equiv \frac{1}{|\Omega|} \int_{\Omega} v(\mathbf{x}) d^3\mathbf{x}.$$

- ▶ Dirichlet-Neumann Boundary Value Problem [3]:

$$\begin{aligned} \Delta v(\mathbf{x}) &= -\frac{1}{D}, \quad \mathbf{x} \in \Omega; \\ \partial_n v(\mathbf{x}) &= 0, \quad \mathbf{x} \in \partial\Omega \setminus \bigcup_j \partial\Omega_{\epsilon_j}; \quad v(\mathbf{x}) = 0, \quad \mathbf{x} \in \bigcup_j \partial\Omega_{\epsilon_j}. \end{aligned}$$



Asymptotic Solutions

The boundary value problem (1) does not admit a known analytic solution. Difficulties arise because of the strongly heterogeneous boundary conditions.

Instead focus on finding high-order asymptotic approximations. Benefits of asymptotic solutions over numerical methods include:

- ▶ Faster computation times.
- ▶ Properties of exact solutions can be extracted.

Asymptotic approximations are of the form

$$v(\mathbf{x}) \sim \epsilon^{-1} v_0(\mathbf{x}) + v_1(\mathbf{x}) + \epsilon \log\left(\frac{\epsilon}{2}\right) v_2(\mathbf{x}) + \epsilon v_3(\mathbf{x}) + \dots$$

where ϵ is the order of magnitude of trap sizes.

Surface Neumann-Green's Function

Of critical importance to the NEP is the surface Neumann-Green's Function, $G_s(\mathbf{x}, \mathbf{x}_j)$, satisfying

$$\begin{aligned} \Delta G_s(\mathbf{x}; \mathbf{x}_j) &= \frac{1}{|\Omega|}, \quad \mathbf{x} \in \Omega; \\ \partial_n G_s(\mathbf{x}; \mathbf{x}_j) &= \delta_s(\mathbf{x} - \mathbf{x}_j), \quad \mathbf{x} \in \partial\Omega; \quad \int_{\Omega} G_s(\mathbf{x}; \mathbf{x}_j) d^3\mathbf{x} = 0. \end{aligned}$$

Using the method of matched asymptotic expansions, the surface Neumann-Green's function appears in the expression for the MFPT as

$$v(\mathbf{x}) = \bar{v} + \sum_{j=1}^N k_j G_s(\mathbf{x}; \mathbf{x}_j), \quad k_j = \text{const.}$$

The Unit Sphere

The special case when Ω is a unit sphere with N holes of radii ϵa_j centred at \mathbf{x}_j respectively yields numerous results [1].

- ▶ Surface Neumann-Green's Function:

$$G_s(\mathbf{x}, \mathbf{x}_j) = \frac{1}{2\pi|\mathbf{x} - \mathbf{x}_j|} + \frac{1}{8\pi} \left(|\mathbf{x}|^2 + 1 \right) + \frac{1}{4\pi} \log \left(\frac{2}{1 - |\mathbf{x}| \cos \gamma + |\mathbf{x} - \mathbf{x}_j|} \right) - \frac{7}{10\pi}$$

- ▶ Mean First Passage Time:

$$v(\mathbf{x}) = \frac{|\Omega|}{2\pi\epsilon DN\bar{c}} \left[1 + \epsilon \log\left(\frac{2}{\epsilon}\right) \sum_{j=1}^N \frac{c_j^2}{2N\bar{c}} - 2\pi\epsilon \sum_{j=1}^N c_j G_s(\mathbf{x}, \mathbf{x}_j) + \frac{2\pi\epsilon}{N\bar{c}} p_c(\mathbf{x}_1, \dots, \mathbf{x}_N) - \frac{\epsilon}{N\bar{c}} \sum_{j=1}^N c_j \kappa_j + \mathcal{O}(\epsilon^2 \log \epsilon) \right].$$

- ▶ Average MFPT:

$$\bar{v} = \frac{|\Omega|}{2\pi\epsilon DN\bar{c}} \left[1 + \epsilon \log\left(\frac{2}{\epsilon}\right) \sum_{j=1}^N \frac{c_j^2}{2N\bar{c}} + \frac{2\pi\epsilon}{N\bar{c}} p_c(\mathbf{x}_1, \dots, \mathbf{x}_N) - \frac{\epsilon}{N\bar{c}} \sum_{j=1}^N c_j \kappa_j + \mathcal{O}(\epsilon^2 \log \epsilon) \right].$$

- ▶ Two important quantities depending only on $\partial\Omega_{\epsilon_j}$:

$$c_j = \frac{2a_j}{\pi} \text{ (trap capacitance)}, \quad \kappa_j = \frac{c_j}{2} \left[2 \log 2 - \frac{3}{2} + \log a_j \right].$$

Self-Interaction Term $p_c(\mathbf{x}_1, \dots, \mathbf{x}_N)$

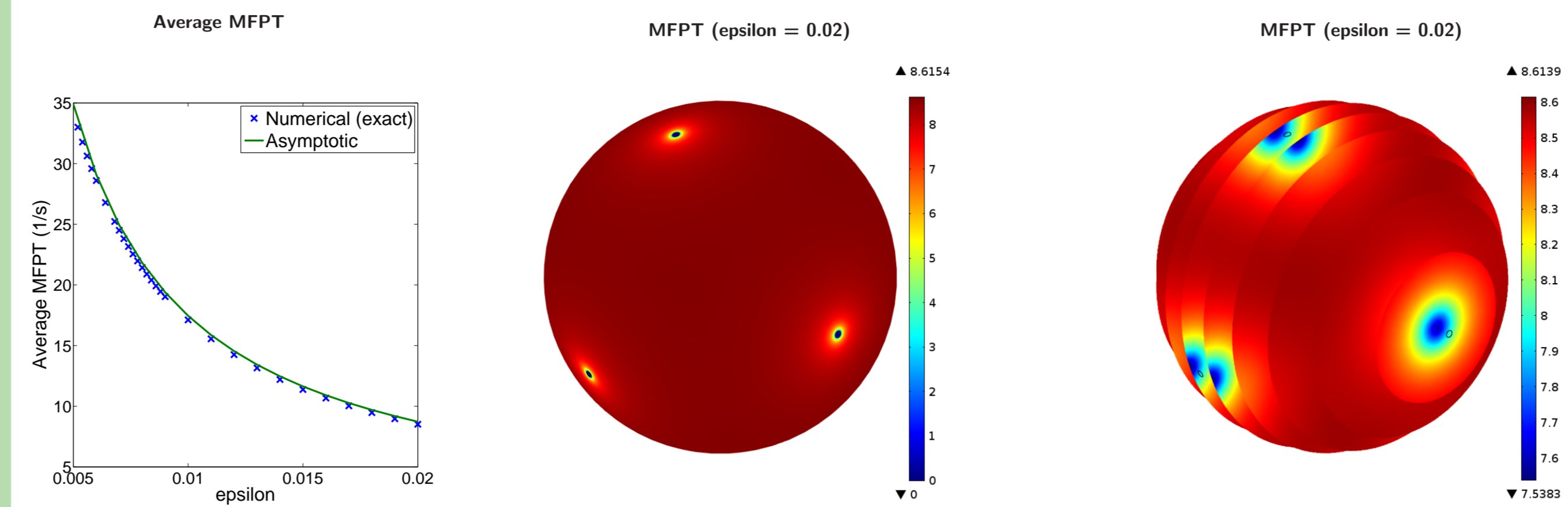
Term $p_c(\mathbf{x}_1, \dots, \mathbf{x}_N)$ appearing in expressions for $v(\mathbf{x})$ and $\bar{v}(\mathbf{x})$ is a self-interaction term.

- ▶ Describes interaction between individual traps \Rightarrow important for optimization.
- ▶ Depends only on $G_s(\mathbf{x}_i, \mathbf{x}_j)$ and each c_j according to

$$p_c(\mathbf{x}_1, \dots, \mathbf{x}_N) = \mathcal{C}^T \mathcal{G}_s \mathcal{C}$$

$$\mathcal{G}_s \equiv \begin{pmatrix} -\frac{9}{20\pi} & G_s(\mathbf{x}_1, \mathbf{x}_2) & \cdots & G_s(\mathbf{x}_1, \mathbf{x}_N) \\ G_s(\mathbf{x}_2, \mathbf{x}_1) & -\frac{9}{20\pi} & \cdots & G_s(\mathbf{x}_2, \mathbf{x}_N) \\ \vdots & \vdots & \ddots & \vdots \\ G_s(\mathbf{x}_N, \mathbf{x}_1) & \cdots & G_s(\mathbf{x}_N, \mathbf{x}_{N-1}) & -\frac{9}{20\pi} \end{pmatrix}, \quad \mathcal{C} \equiv \begin{pmatrix} c_1 \\ \vdots \\ c_N \end{pmatrix}.$$

Asymptotic and Numerical MFPT for Unit Sphere with Six Identical Traps



Singer, Schuss, and Holcman Approximation

When Ω is a general three-dimensional domain, previous results are limited to the case of one trap $\partial\Omega_{\epsilon}$ of radius ϵ (i.e. $a = 1$) located at \mathbf{x}_0 [4].

- ▶ Surface-Neumann Green's Function:

$$G_s(\mathbf{x}, \mathbf{x}_0) = \frac{1}{2\pi|\mathbf{x} - \mathbf{x}_j|} - \frac{H(\mathbf{x}_0)}{4\pi} \log |\mathbf{x} - \mathbf{x}_0| + v_s(\mathbf{x}, \mathbf{x}_0),$$

where $v_s(\mathbf{x}, \mathbf{x}_j)$ is an unknown bounded function of $\mathbf{x}, \mathbf{x}_j \in \Omega$.

- ▶ Average MFPT:

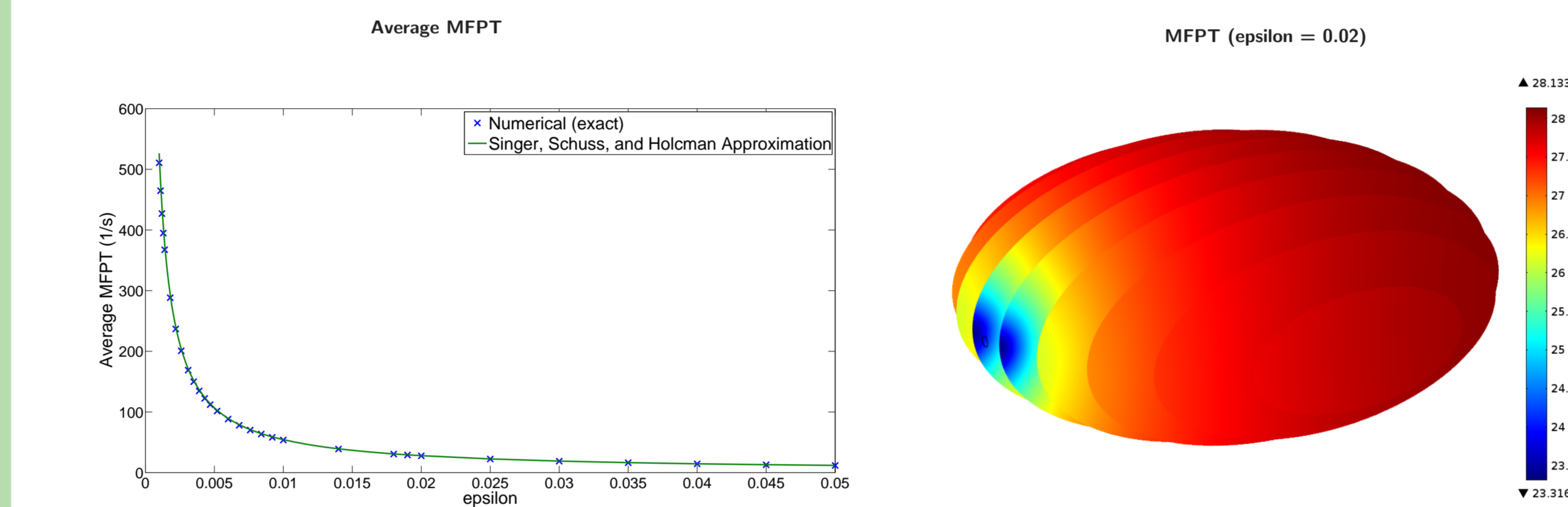
$$\bar{v} \equiv \frac{|\Omega|}{4\epsilon D} \left[1 + \frac{H(\mathbf{x}_0)}{\pi} \epsilon \log \epsilon + \mathcal{O}(\epsilon) \right]^{-1}.$$

Limitations of this approach are:

- ▶ Approximation is only valid for one absorbing window.
- ▶ No asymptotic expression for the (non-averaged) MFPT is given.
- ▶ Error bound of $\mathcal{O}(\epsilon)$ is worse than that for sphere.

Singer, Schuss, and Holcman Approximation for Oblate Spheroid with One Trap

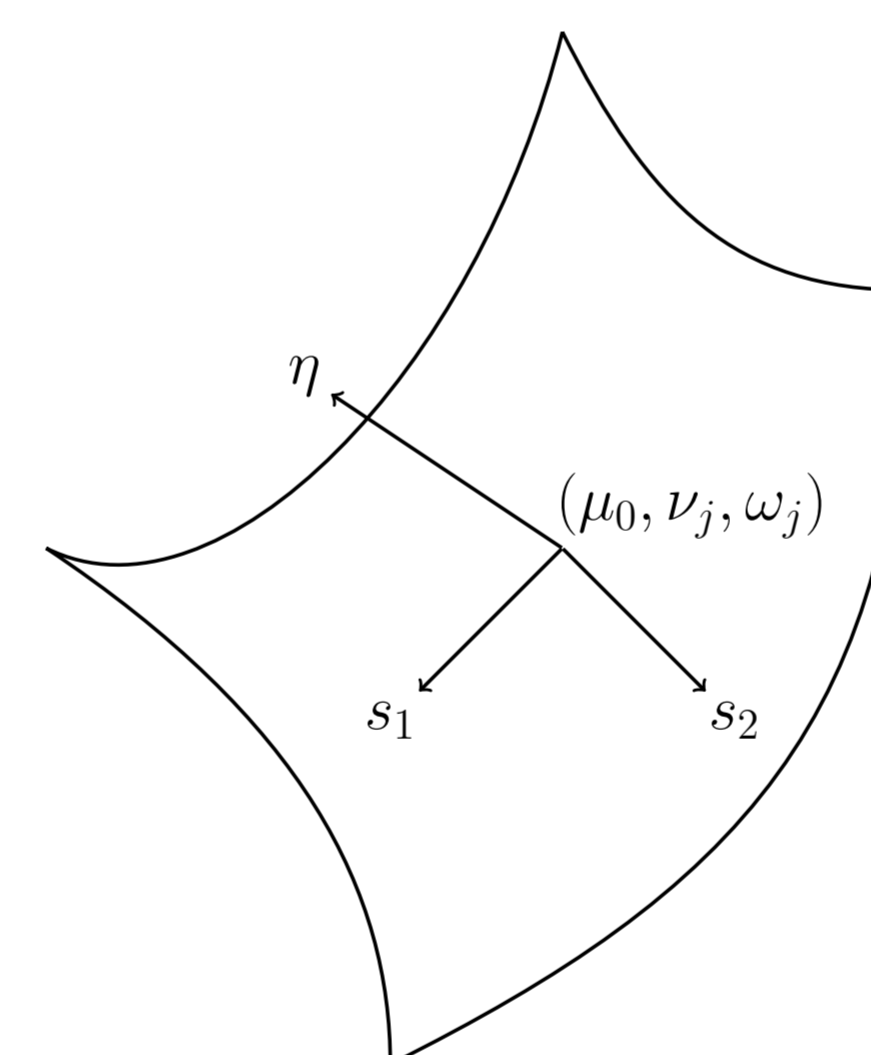
Oblate Spheroid with One Trap



Towards a Wider Class of Three-Dimensional Domains

Consider a class of 3D domains where boundary is a coordinate surface for some orthonormal coordinate system (μ, ν, ω) . Then assume the coordinate surface is $\mu = \mu_0$.

- ▶ Examples: spheres, spheroids, ellipsoids, surfaces of rotation.
- ▶ N traps located at (μ_0, ν_j, ω_j) for $j = 1, \dots, N$.
- ▶ h_μ, h_ν, h_ω : scale factors of particular coordinate system.
- ▶ Local stretched coordinates:



$$\eta = -h_\mu \frac{\mu - \mu_0}{\epsilon}, \quad s_1 = h_\nu \frac{\nu - \nu_j}{\epsilon}, \quad s_2 = h_\omega \frac{\omega - \omega_j}{\epsilon}.$$

Local Form of Surface Neumann-Green's Function

Using the expression for the surface Neumann-Green's function (1) and introducing the local stretched coordinates gives:

$$G_s(\eta, s_1, s_2; \mathbf{x}_j) = \frac{1}{2\pi\rho\epsilon} - \frac{H(\mathbf{x}_j)}{4\pi} \log \frac{\rho}{2} + g_0(\eta, s_1, s_2; \mathbf{x}_j) + \epsilon \log \frac{\rho}{2} g_1(\eta, s_1, s_2; \mathbf{x}_j) + \mathcal{O}(\epsilon),$$

where $\rho = \sqrt{\eta^2 + s_1^2 + s_2^2}$ and g_0 and g_1 are bounded functions depending on the geometry at \mathbf{x}_j .

Method of Matched Asymptotic Expansions

The solution is formulated in terms of inner and outer solutions, each satisfying a corresponding problem.

Inner Problem (near \mathbf{x}_j)

- ▶ Local stretched coordinates (η, s_1, s_2) .

$$w(\eta, s_1, s_2) \sim \frac{1}{\epsilon} w_0 + \log \frac{\rho}{2} w_1 + w_2 + \mathcal{O}(\epsilon).$$

- ▶ Domain: $\eta \geq 0, s_1, s_2 \in \mathbb{R}$.

- ▶ Linear PDE: $\Delta_{(\eta, s_1, s_2)} w_k = \delta_{k2} \mathcal{L} w_0$.

(\mathcal{L} is a second-order linear differential operator.)

- ▶ Boundary Conditions:

$$\begin{aligned} \partial_\eta w_k &= 0, & \eta &= 0, & s_1^2 + s_2^2 &\geq a_j^2, \\ w_k &= 0, & \eta &= 0, & s_1^2 + s_2^2 &\leq a_j^2. \end{aligned}$$

Outer Problem (far from \mathbf{x}_j)

- ▶ Global coordinates (μ, ν, ω) .

$$v(\mu, \nu, \omega) \sim \frac{1}{\epsilon} v_0 + v_1 + \epsilon \log \frac{\rho}{2} v_2 + \epsilon v_3 + \mathcal{O}(\epsilon).$$

- ▶ Domain: $(\mu, \nu, \omega) \in \Omega$.

- ▶ PDE: $\Delta v_k = -\frac{1}{D} \delta_{k1}$.

- ▶ Boundary Conditions:

$$\partial_n v_k = 0, \quad \mathbf{x} \in \partial\Omega \setminus \{\mathbf{x}_1, \dots, \mathbf{x}_N\}.$$

Matched Asymptotic Expansions Condition

$$\frac{1}{\epsilon} w_0 + \log \frac{\rho}{2} w_1 + w_2 + \dots \sim \frac{1}{\epsilon} v_0 + v_1 + \epsilon \log \frac{\rho}{2} v_2 + \epsilon v_3 + \dots$$

- ▶ Collect like coefficients of ϵ and sequentially solve for w_k and v_k using the inner problem, the outer problem, and the matching condition.

Proposed Asymptotic Solutions for MFPT and Average MFPT

- ▶ Assumptions: $g_1 = 0$ and $w_2 \sim \frac{v_0 b_j}{\rho}$.

- ▶ MFPT:

$$v(\mathbf{x}) = \frac{|\Omega|}{2\pi\epsilon DN\bar{c}} \left[1 - \epsilon \log\left(\frac{\epsilon}{2}\right) \frac{\sum_{j=1}^N c_j^2 H(\mathbf{x}_j)}{2N\bar{c}} - 2\pi\epsilon \sum_{j=1}^N c_j G_s(\mathbf{x}, \mathbf{x}_j) + \frac{2\pi\epsilon}{N\bar{c}} p_c(\mathbf{x}_1, \dots, \mathbf{x}_N) - \frac{\epsilon}{N\bar{c}} \sum_{j=1}^N b_j + \mathcal{O}(\epsilon^2 \log \epsilon) \right].$$

- ▶ Average MFPT:

$$\bar{v} = \frac{|\Omega|}{2\pi\epsilon DN\bar{c}} \left[1 - \epsilon \log\left(\frac{\epsilon}{2}\right) \frac{\sum_{j=1}^N c_j^2 H(\mathbf{x}_j)}{2N\bar{c}} + \frac{2\pi\epsilon}{N\bar{c}} p_c(\mathbf{x}_1, \dots, \mathbf{x}_N) - \frac{\epsilon}{N\bar{c}} \sum_{j=1}^N b_j + \mathcal{O}(\epsilon^2 \log \epsilon) \right].$$

- ▶ $\tilde{p}_c(\mathbf{x}_1, \dots, \mathbf{x}_N)$ is a modified version of $p_c(\mathbf{x}_1, \dots, \mathbf{x}_N)$ for the sphere, depending only on $G_s(\mathbf{x}_i, \mathbf{x}_j)$.
- ▶ b_j is modified version of κ_j for the sphere, determined by the far field behaviour of w_2 .

Testing Procedure and COMSOL

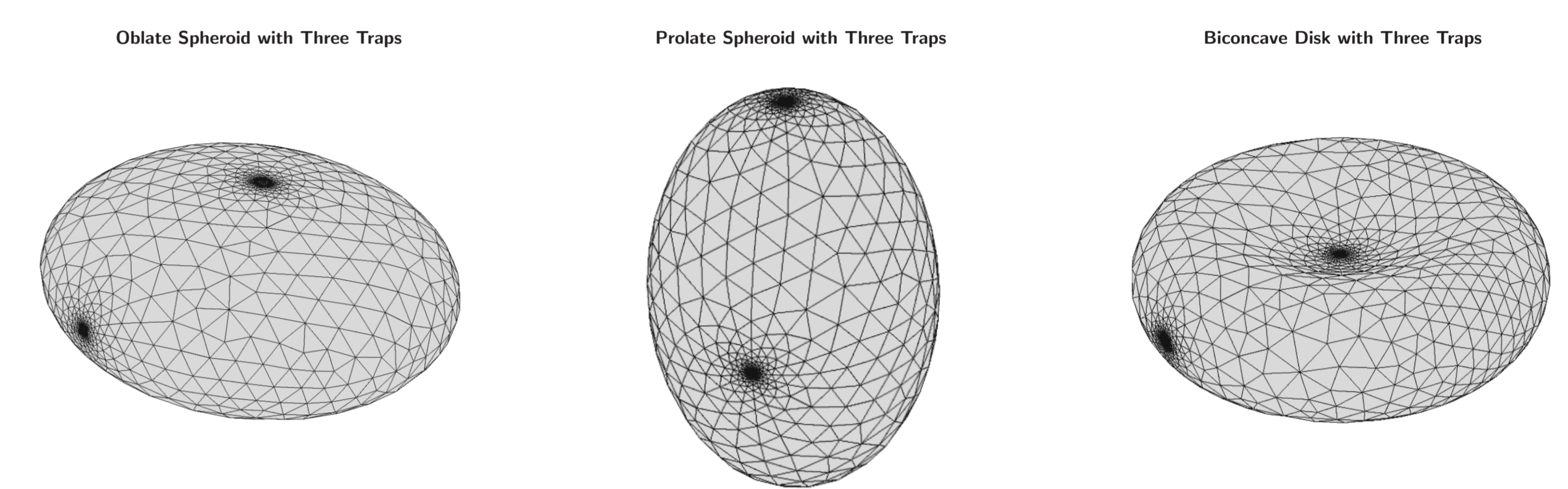
Used oblate spheroid, prolate spheroid, and biconcave disk geometries.

- ▶ Provide range of local curvatures.
- ▶ Represent different biological cells.

COMSOL Multiphysics 4.3b software used for numerical results.

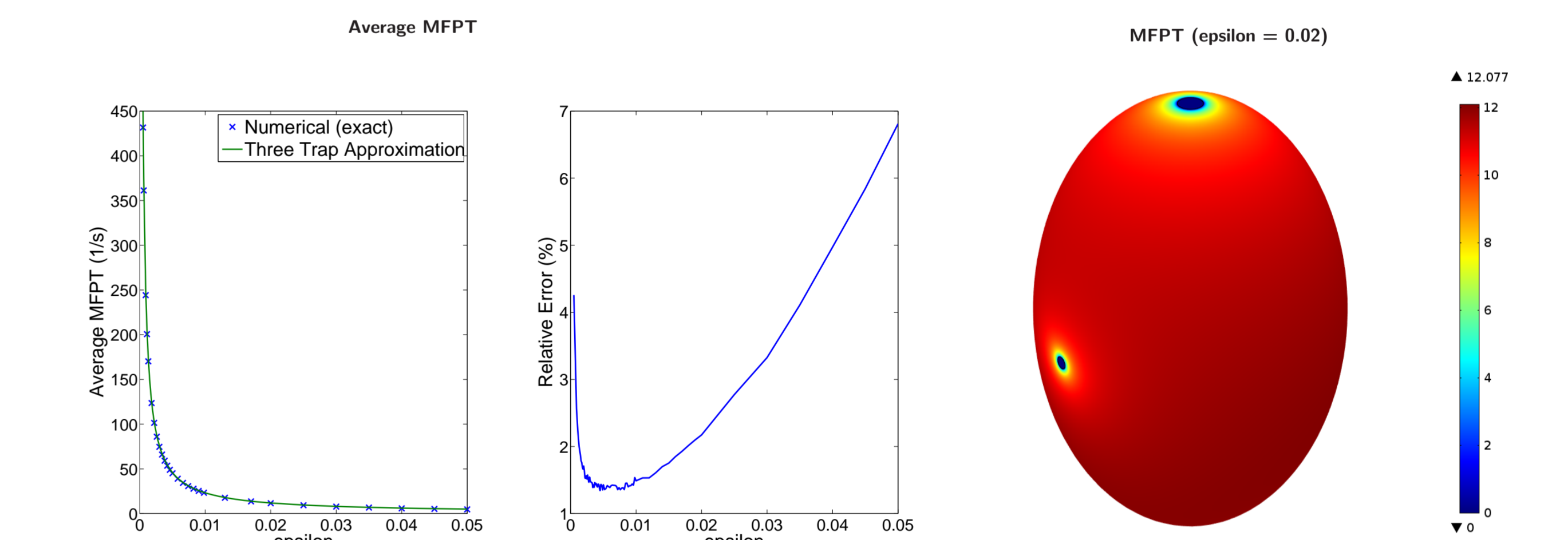
- ▶ Finite element PDE solver.
- ▶ Tetrahedral mesh.

Numerical results for two and three traps of equal and different sizes compared to proposed multi-trap approximation in MATLAB.

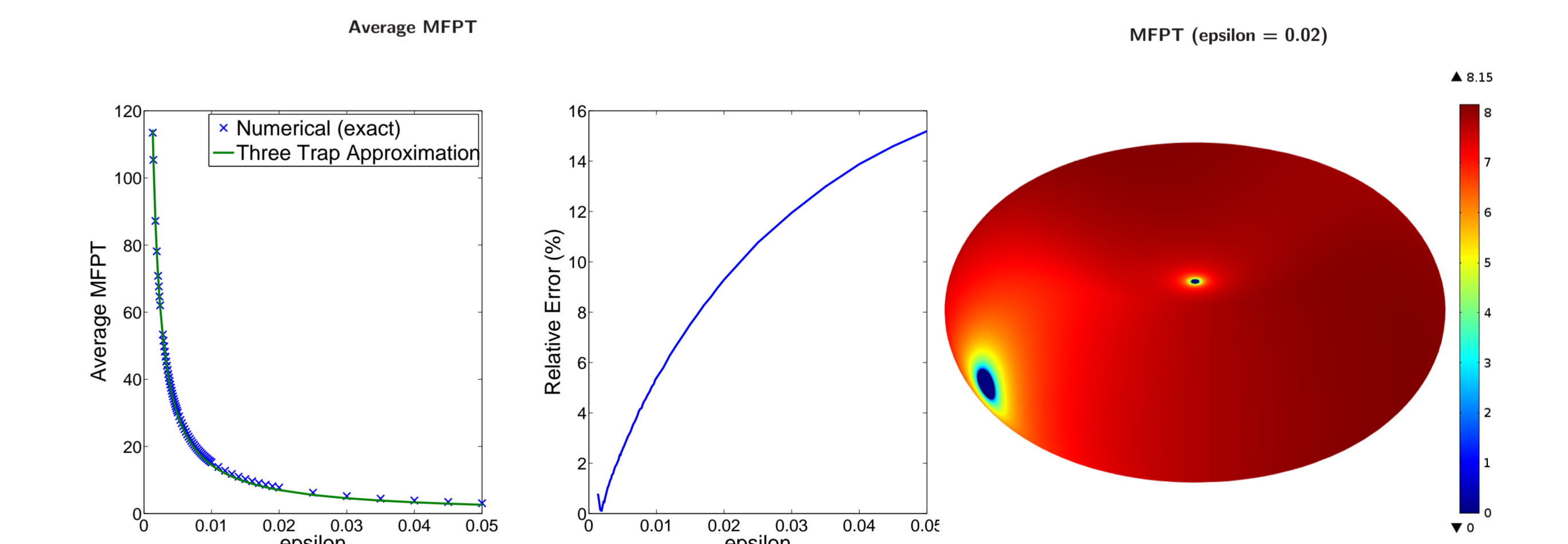


Results for Three Traps of Different Sizes

Prolate Spheroid



Biconcave Disk



Conclusions

- ▶ Expressions for the MFPT and average MFPT were developed for a more general class of three dimensional domains.
- ▶ The average MFPT values following from the proposed asymptotic formulae were found to be in close agreement with numerical simulation results.

Future Research

- ▶ Comparison to numerical simulation for a more extensive variety of geometries.
- ▶ Rigorous justification of assumptions used for proposed MFPT and average MFPT formulae.
- ▶ Study of dilute trap limit of homogenization theory for non-spherical domains [2].

References

- [1] Alexei F. Cheviakov, Michael J. Ward, and Ronny Straube. An asymptotic analysis of the mean first passage time for narrow escape problems. II. The sphere. *Multiscale Model. Simul.*, 8(3):836–870, 2010.
- [2] Cyrill B. Muratov and Stanislav Y. Shvartsman. Boundary homogenization for periodic arrays of absorbers. *Multiscale Model. Simul.*, 7(1):44–61, 2008.
- [3] Zeev Schuss. *Theory and applications of stochastic differential equations*. John Wiley & Sons Inc., New York, 1980. Wiley Series in Probability and Statistics.
- [4] A. Singer, Z. Schuss, and D. Holcman. Narrow escape and leakage of Brownian particles. *Phys. Rev. E* (3), 78(5):051111, 8, 2008.