

The Narrow Escape Problem

Consider a particle that undergoes **Brownian motion** while confined to the interior of a domain Ω . The boundary of this domain is made up of almost entirely a **reflecting portion**, $\partial\Omega_r$ and a relatively small **absorbing portion**, $\partial\Omega_a$. Define $v(\mathbf{x}), \mathbf{x} \in \Omega$ as the expectation time a particle will stay in the domain starting at point \mathbf{x} .

Dirichlet-Neumann Boundary Problem

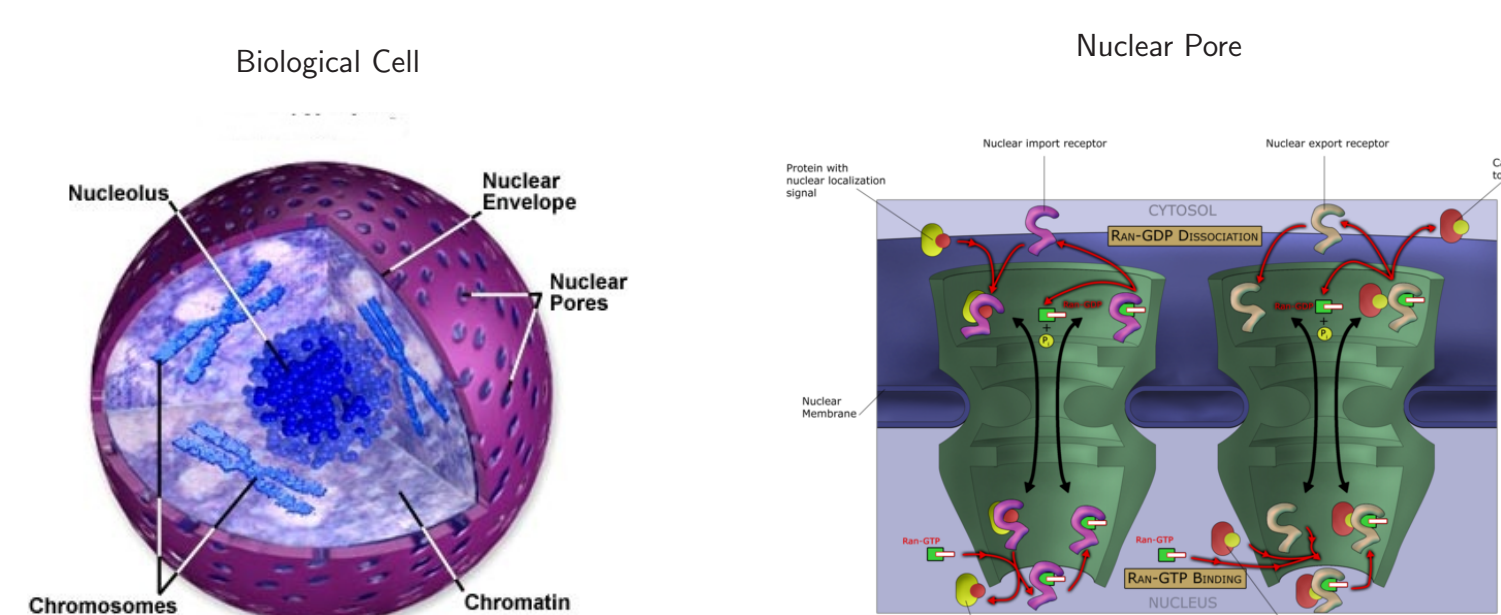
$$\Delta v = -\frac{1}{D}, \mathbf{x} \in \Omega \quad v = 0, \mathbf{x} \in \partial\Omega_a \quad \partial_r v = 0, \mathbf{x} \in \partial\Omega_r, \quad (1)$$

where D is the diffusion constant for the system.

Mean First Passage Time (MFPT):

$$\bar{v} = \int v(\mathbf{x}) d\mathbf{x} \quad (2)$$

Applications



- Nuclear export of messenger RNA through nuclear pores
- Reciprocal of MFPT acts as a first order rate constant

The Asymptotic Solution For the Unit Sphere

An **asymptotic solution** was derived for the unit sphere with traps located at \mathbf{x}_i using the method of **matched asymptotic expansions**.

Assumptions:

- Small trap sizes, $\epsilon \ll 1$
- Well separated traps, $|\mathbf{x}_i - \mathbf{x}_j| \ll \epsilon$

$$\bar{v} = \frac{|\Omega|}{4\epsilon DN} \left[1 + \frac{\epsilon}{\pi} \log\left(\frac{2}{\epsilon}\right) + \frac{\epsilon}{\pi} \left(-\frac{9N}{5} + 2(N-2)\log 2 + \frac{3}{2} + \frac{4}{N} \mathcal{H}(x_1, \dots, x_N) \right) \right] \quad (3)$$

$\mathcal{H}(x_1, \dots, x_N)$ is the **interaction energy** defined by

$$\mathcal{H}(x_1, \dots, x_N) = \sum_{i=1}^N \sum_{j=i+1}^N \left(\frac{1}{|\mathbf{x}_i - \mathbf{x}_j|} - \frac{1}{2} \log |\mathbf{x}_i - \mathbf{x}_j| - \frac{1}{2} \log(2 + |\mathbf{x}_i - \mathbf{x}_j|) \right) \quad (4)$$

More generally, for unequal trap lengths the MFPT was found to be

$$\bar{v} = \frac{|\Omega|}{2\pi\epsilon DN\bar{c}} \left[1 + \epsilon \log\left(\frac{2}{\epsilon}\right) \frac{\sum_{j=1}^N c_j^2}{2N\bar{c}} + \frac{2\pi\epsilon}{N\bar{c}} \mathcal{G}^T \mathcal{G} \mathcal{C} - \frac{\epsilon}{N\bar{c}} \sum_{j=1}^N c_j \kappa_j \right], \quad (5)$$

where c_j are the elements of the capacitance vector \mathcal{C} , given by $c_j = 2a_j/\pi$ where $a_j\epsilon$ is the radius of the trap corresponding to index j . \bar{c} is the mean of the capacitance vector, and \mathcal{G} is the **surface Green's function matrix** given by

$$\mathcal{G}_{ij} = G_s(x_i; x_j) = \frac{1}{2\pi|\mathbf{x}_i - \mathbf{x}_j|} + \frac{1}{4\pi} \log\left(\frac{2}{1 - \cos\gamma + |\mathbf{x}_i - \mathbf{x}_j|}\right) - \frac{9}{20\pi}, \quad i \neq j, \quad (6)$$

$$\mathcal{G}_{ii} = \frac{-9}{20\pi},$$

where γ is the angle between traps \mathbf{x}_i and \mathbf{x}_j .

The Homogenization Problem

The **homogenization problem** deals with the limit of many equally spaced traps that cover a fixed percentage, σ , of the boundary. In this limit, it is possible to replace the Dirichlet-Neumann boundary conditions with an equivalent

Robin Boundary Condition Problem:

$$\Delta v_H = -\frac{1}{D}, \quad v_H \in \Omega; \quad \epsilon \partial_r v_H + \kappa v_H = 0, \quad v \in \partial\Omega, \quad (7)$$

where κ is a factor that depends on the geometry of the domain. This equation can be solved immediately for the sphere; from (2), one gets

$$\bar{v}_H = \frac{1}{15D} + \frac{1}{3D\kappa}. \quad (8)$$

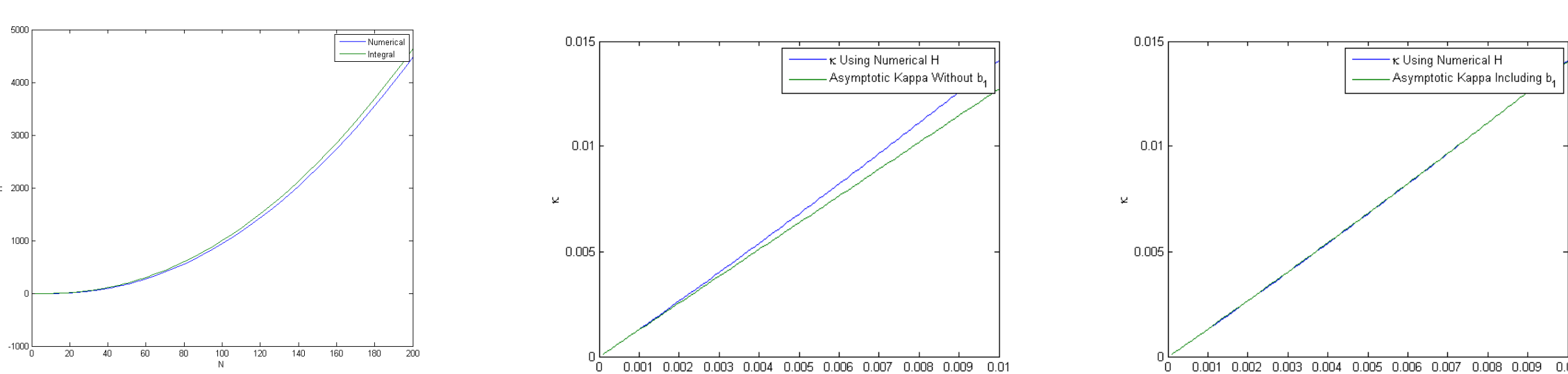
► \mathcal{H} can be estimated by replacing the sum with an integral over the sphere. One gets

$$\bar{v} \approx \frac{|\Omega|}{4\epsilon DN} \left[1 - \frac{\epsilon}{\pi} \log \epsilon + \frac{\epsilon N}{\pi} \left(\frac{1}{5} + \frac{4b_1}{\sqrt{N}} \right) \right]. \quad (9)$$

Comparison of Numerical \mathcal{H} to Integral \mathcal{H}

κ without b_1

κ with fitted b_1



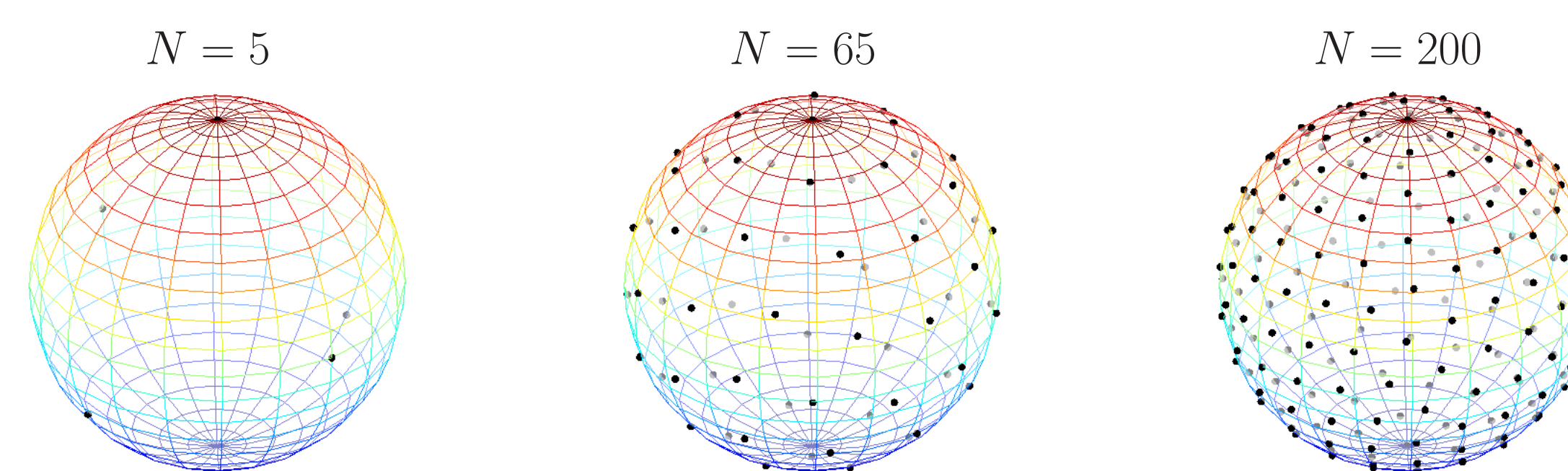
- Expect leading order term is correct.
- Consider leading order terms $\mathcal{H} \approx N^2 \frac{1}{2}(1 - \log 2) + b_1 N^{3/2}$ where b_1 is a constant.
- Equate \bar{v} and \bar{v}_H in the limit $N \rightarrow \infty$ to estimate κ :

$$\kappa = \frac{4\sigma}{\pi + 8b_1\sqrt{\sigma}}. \quad (10)$$

Fitting κ to the full numerical results for \mathcal{H} gives us $b_1 = -0.3672$, because of the reasonable agreement shown above (3) provides a quick way to estimated \bar{v} in the high N , low σ limit.

Optimal Trap Locations on the Unit Sphere

The optimal trap configuration corresponds to the **minimized MFPT**, or **maximized diffusion rate**. Examples of optimal configurations:



Using **global optimization software** (GANSO, LGO), results were computed up to $N = 200$. The computations take a long time since this is minimizing a function in $(2N - 3)$ -dimensional space.

The Topological Derivative

In [3], the notion of a **topological derivative** is developed. Consider a **functional** $\mathcal{J} : \Omega \mapsto \mathbb{R}$; define the topological derivative as

$$\mathfrak{T}(\mathbf{x}) = \lim_{\rho \rightarrow 0} \frac{\mathcal{J}(\Omega \setminus B_\rho(\mathbf{x})) - \mathcal{J}(\Omega)}{|B_\rho(\mathbf{x})|}, \quad (11)$$

where $B_\rho(\mathbf{x})$ is a ball of radius ρ located at point $\mathbf{x} \in \Omega$.

► Example:



Topological Derivative for Narrow Escape on Unit Sphere

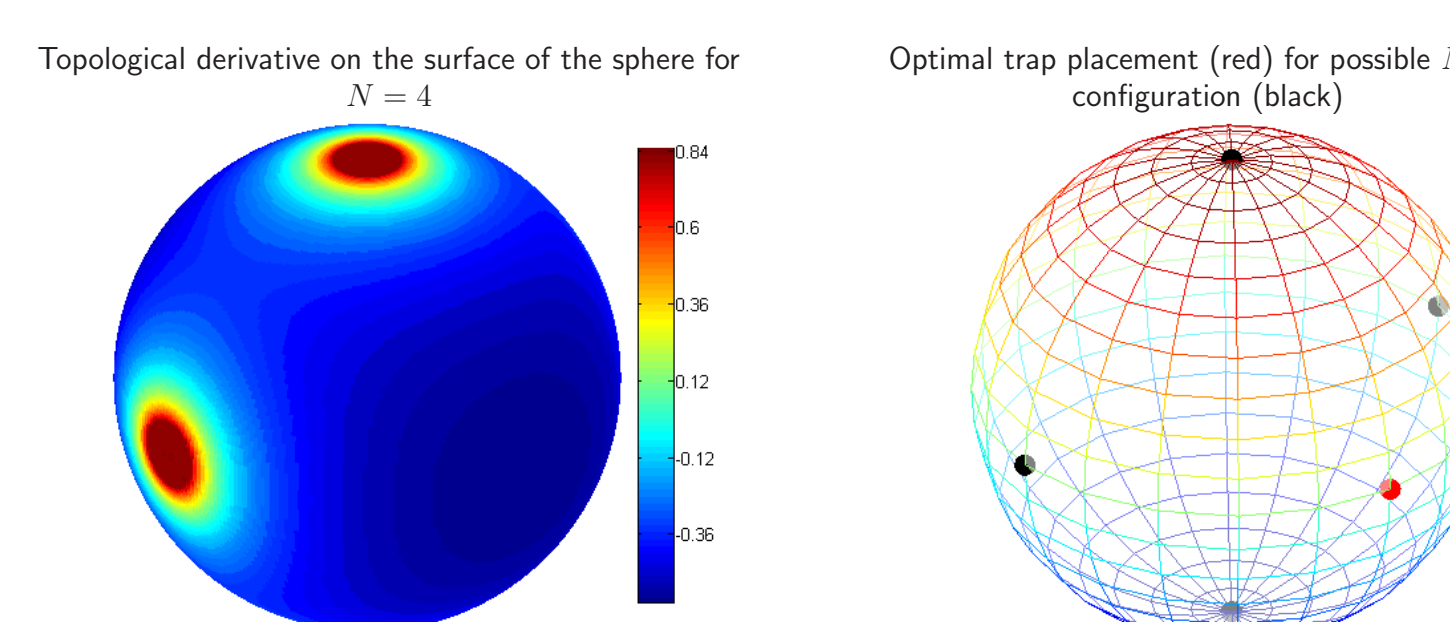
In our case the shape functional \mathcal{J} is \bar{v} , the MFPT and $B_\rho(\mathbf{x})$ is a trap of radius ρ centered at point \mathbf{x} on the sphere.

► Topological derivative calculated using (5)

$$\mathfrak{T}(\mathbf{x}) = \frac{|\Omega|}{2\pi\epsilon DN\bar{c}} \left[-\log\left(\frac{2}{\epsilon}\right) \left(\frac{1}{\pi N^2} \right) + \frac{3-4\log 2}{2\pi N^2} + \frac{8N \sum_{j=1}^N \mathcal{G}_{j(N+1)} - 4 \sum_{i,j} \mathcal{G}_{ij}}{N^2} \right] \quad (12)$$

► Optimal trap location is when $\mathfrak{T}(\mathbf{x})$ is minimized

► Topological derivative is minimized when $\sum_{j=1}^N \mathcal{G}_{j(N+1)}$ is minimized



We iterate this process to obtain a **rough estimate** for the minimum energy.

Application of Topological Derivative to Computation of Optimal Arrangements

For two particles, the interaction energy is a strictly decreasing function in terms of the separation distance $|\mathbf{x}_i - \mathbf{x}_j|$. We introduce the pseudo-force vector

$$\vec{F}_{ij} = \left(\frac{1}{|\mathbf{x}_i - \mathbf{x}_j|} - \log(|\mathbf{x}_i - \mathbf{x}_j|) - \log(2 + |\mathbf{x}_i - \mathbf{x}_j|) \right) \mathbf{e}_{ij} \quad (13)$$

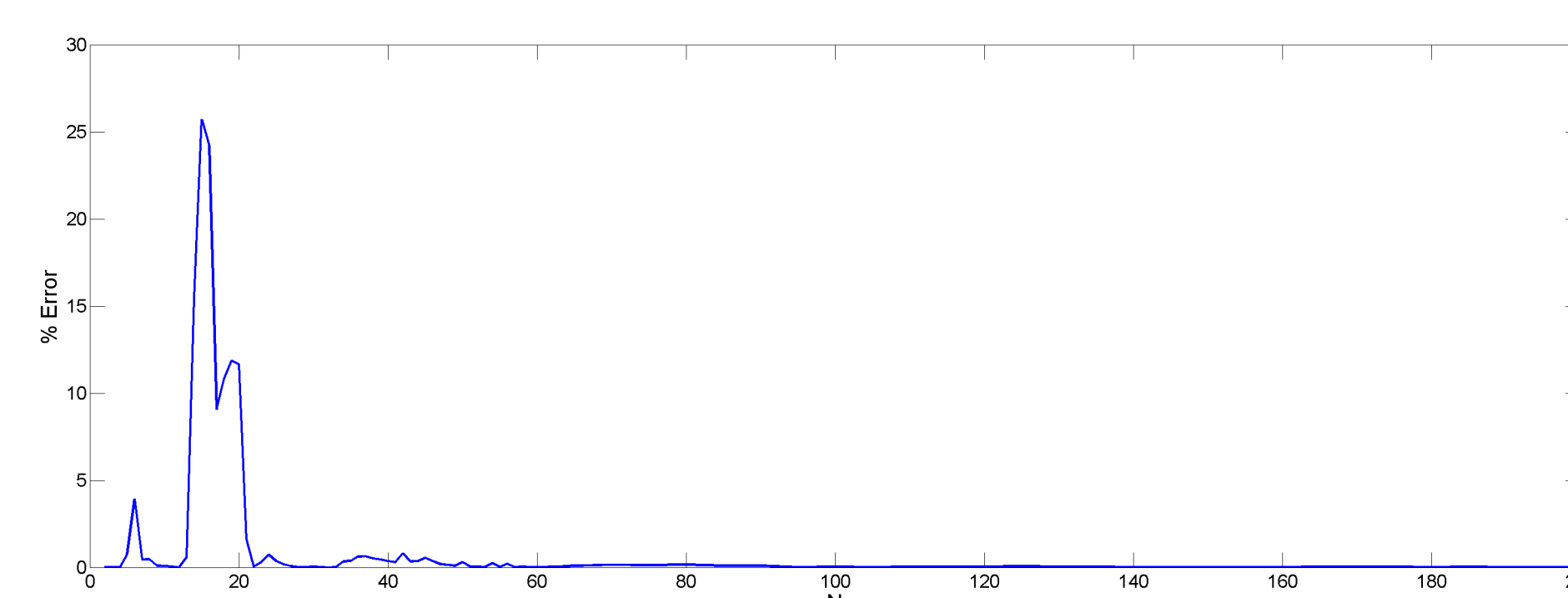
And the total force on a trap \mathbf{x}_i is thus

$$\vec{F}_i = \sum_{j=1, j \neq i}^N \vec{F}_{ij} \quad (14)$$

We start by introducing a new particle in the point of the minimal topological derivative. Then we move each particle a distance proportional to \vec{F}_i . Each particle is subsequently pushed back to the sphere in the normal direction. This process is iterated until a **local energy minimum** is achieved, which may or may not be the **global energy minimum**.

Testing the Simulation against Known Results

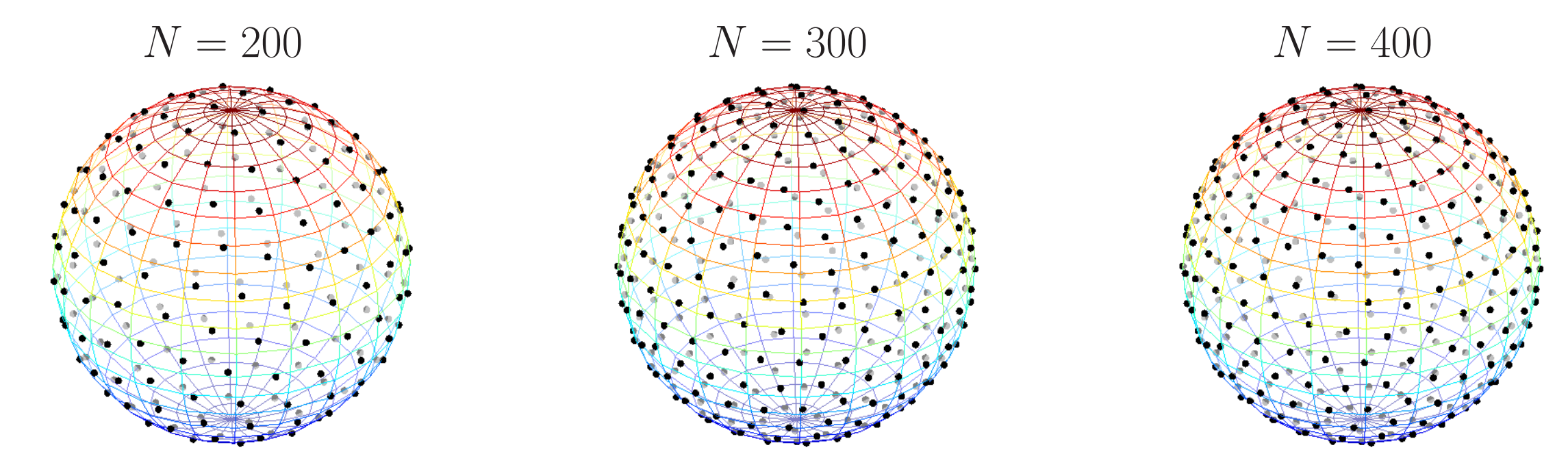
Results for the values of \mathcal{H} from the above simulation were compared against the numerical global optimization results:



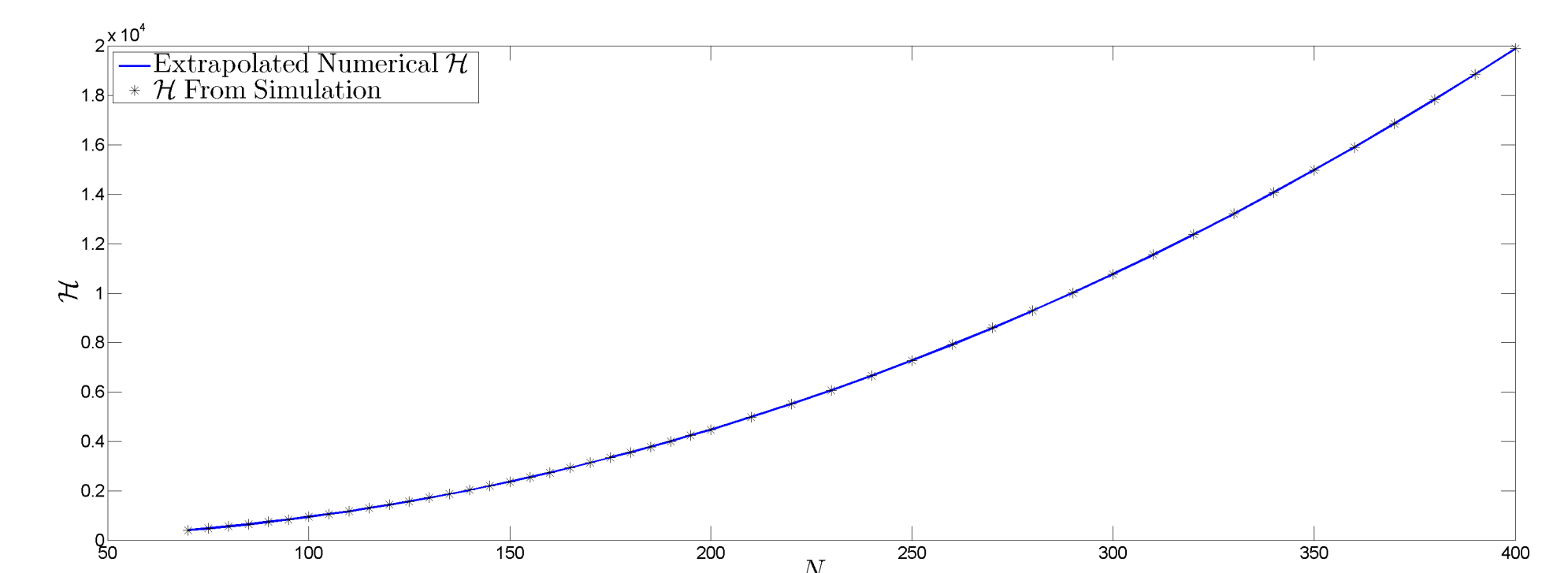
The topological derivative-based simulation showed to be accurate within 0.01% for $N \geq 100$. The simulation is naturally much faster compared to full global optimization.

Simulation Results

Optimal trap locations were computed up to $N = 400$ in steps of 10.



The interaction energies from the simulation results were compared against a **polynomial extrapolation** of the $N = 2..200$ numerical results up to $N=400$:



Simulation \mathcal{H} matches extrapolated numerical \mathcal{H} within **0.005%**.

N^2 Conjecture

Asymptotic Motivation:

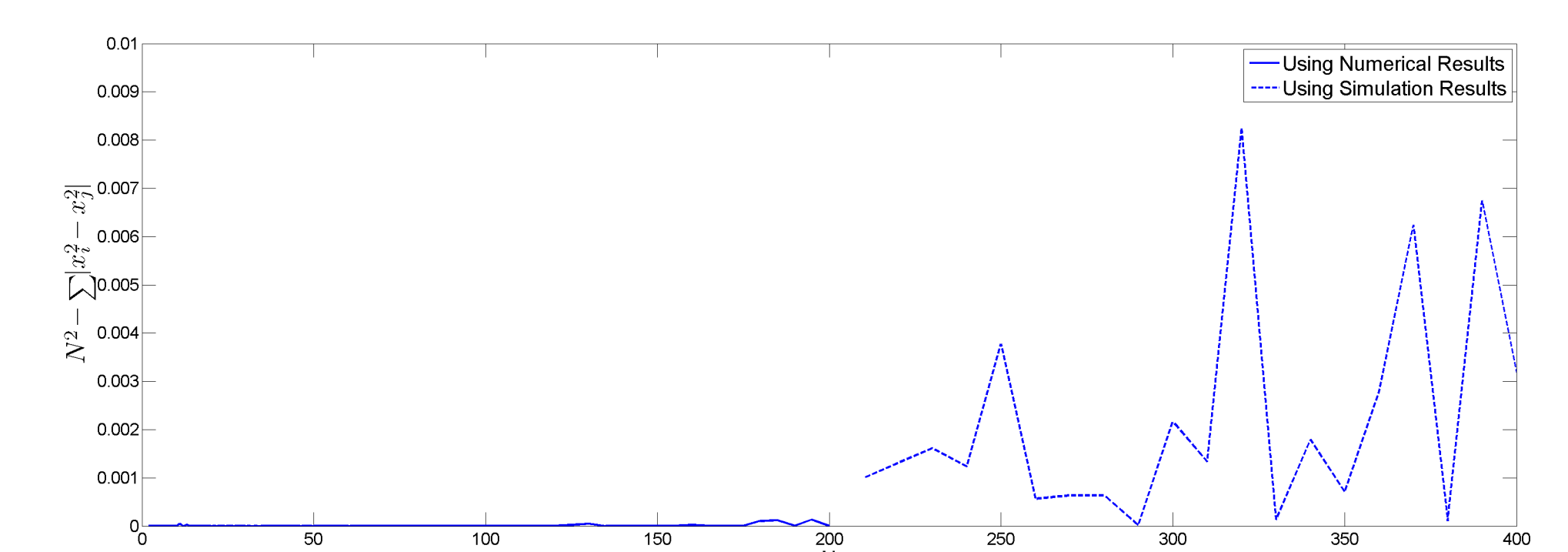
- In the limit of **many traps** we can **approximate** the sum of the trap square distances as an **integral**

$$\sum |\mathbf{x}_i - \mathbf{x}_j|^2 \approx \frac{N^2}{8\pi} \int_0^\pi \int_0^{2\pi} (2 - 2\cos\theta) \cdot \sin\theta d\phi d\theta = N^2. \quad (15)$$

We conjecture that for an arrangement of traps in the minimal energy configuration,

$$N^2 = \sum |\mathbf{x}_i - \mathbf{x}_j|^2 \quad (16)$$

The difference between N^2 and numerical/simulation \mathcal{H} was plotted up to $N = 400$:



Since the **absolute differences were less than 0.01** for all computed N , we conclude that the conjecture holds at least for $N \leq 400$.

Conclusions

- An estimate of κ was obtained for the unit sphere.
- Simulation results shown to **match previous results within 0.01%** for the $N \geq 100$ region.
- Topological derivative-based Simulation provides a **faster** method of computing the optimal trap arrangement on the sphere.
- $N^2 - \sum |\mathbf{x}_i - \mathbf{x}_j|^2 < 0.01$ for all computed N .

Further Research Directions / Open Problems

- Rigorous derivation of κ for the homogenization limit.
- Simulation could be programmed with higher precision.
- Computing the topological derivative without the use of the asymptotic formula for \mathcal{H} .
- Rigorous justification for N^2 conjecture.
- Asymptotic solution of the narrow escape problem for on an arbitrary domain $\Omega \in \mathbb{R}^3$.

References

- [1] M. J. Ward A. F. Cheviakov and R. Straube. An Asymptotic Analysis of the Mean First Passage Time for Narrow Escape Problems: Part II: The Sphere. *SIAM Multiscale Modeling and Simulation*, 2009.
- [2] A. Peirce M. J. Ward, S. Pillay and T. Kolokolnikov. An Asymptotic Analysis of the Mean First Passage Time for Narrow Escape Problems: Part I: Two-Dimensional Domains. *SIAM Multiscale Modeling and Simulation*, Vol. 8, No. 3, pp. 803-835., 2010.
- [3] J. Sokolowski and A. Zochowski. On the Topological Derivative in Shape Optimization. *SIAM Journal on Control and Optimization*, 1997.