

Conservation Laws of Fluid Dynamics Models

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- 1 Fluid Dynamics Equations
- 2 CLs of Constant-Density Euler and N-S Equations
- 3 CLs of Helically Invariant Flows
- 4 CLs of An Inviscid Model in Gas Dynamics
- 5 CLs of a Surfactant Flow Model
- 6 Discussion

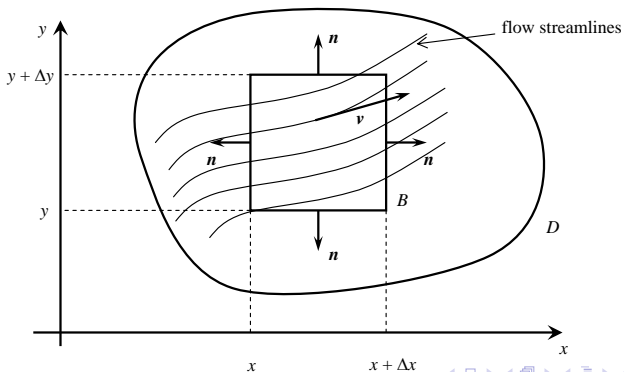
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Fluid/gas flow in 3D

- Independent variables: t, x, y, z .
- Dependent variables: $\mathbf{u} = (u^1, u^2, u^3) = (u, v, w)$; p ; ρ .

Fluid/gas flow in 3D

- Independent variables: t, x, y, z .
- Dependent variables: $\mathbf{u} = (u^1, u^2, u^3) = (u, v, w)$; p ; ρ .
- 2D picture:



- Euler equations:

$$\rho_t + \nabla \cdot (\rho \mathbf{u}) = 0,$$

$$\rho(\mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u}) + \nabla p = 0.$$

- Navier-Stokes equations (viscosity $\nu = \text{const}$):

$$\rho_t + \nabla \cdot (\rho \mathbf{u}) = 0,$$

$$\rho(\mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u}) + \nabla p - \nu \nabla^2 \mathbf{u} = 0.$$

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- 4 equations, 5 unknowns. Closure required.

- Euler equations:

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$$\rho(\mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u}) + \nabla p - \nu \nabla^2 \mathbf{u} = 0.$$

- Closure e.g. 1, homogeneous flow (e.g., water):

$$\rho = \text{const}, \quad \text{div } \mathbf{u} = 0.$$

- Euler equations:

$$\rho_t + \nabla \cdot (\rho \mathbf{u}) = 0,$$

$$\rho(\mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u}) + \nabla p = 0.$$

- Navier-Stokes equations (viscosity $\nu = \text{const}$):

$$\rho_t + \nabla \cdot (\rho \mathbf{u}) = 0,$$

$$\rho(\mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u}) + \nabla p - \nu \nabla^2 \mathbf{u} = 0.$$

- Closure e.g. 2, incompressible flow:

$$\text{div } \mathbf{u} = 0,$$

$$\rho_t + \mathbf{u} \cdot \nabla \rho = 0.$$

- Euler equations:

$$\rho_t + \nabla \cdot (\rho \mathbf{u}) = 0,$$

$$\rho(\mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u}) + \nabla p = 0.$$

- Navier-Stokes equations (viscosity $\nu = \text{const}$):

$$\rho_t + \nabla \cdot (\rho \mathbf{u}) = 0,$$

$$\rho(\mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u}) + \nabla p - \nu \nabla^2 \mathbf{u} = 0.$$

- Other closure choices: ideal gas/adiabatic, isothermal, polytropic (gas dynamics), etc...

- Euler equations:

$$\rho_t + \nabla \cdot (\rho \mathbf{u}) = 0,$$

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- Navier-Stokes equations (viscosity $\nu = \text{const}$):

$$\rho_t + \nabla \cdot (\rho \mathbf{u}) = 0,$$

$$\rho(\mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u}) + \nabla p - \nu \nabla^2 \mathbf{u} = 0.$$

- Multiple other fluid models exist.

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Constant-density Euler equations:

$$\nabla \cdot \mathbf{u} = 0,$$

$$\mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u} + \nabla p = 0.$$



Constant-density Euler equations:

$$\nabla \cdot \mathbf{u} = 0,$$

$$\mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u} + \nabla p = 0.$$

- CLs in a general setting.
- Additional CLs in a symmetric setting (e.g., axisymmetric).
- More additional CLs in a reduced setting (e.g., planar flow).

Euler equations:

$$\nabla \cdot \mathbf{u} = 0,$$

$$\mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u} + \nabla p = 0.$$

- Some conservation laws known “forever”, e.g., [Batchelor (2000)].
- Kovalevskaya form w.r.t. x, y, z .
- It remains an **open problem** to determine the upper bound of the CL order for the Euler system.
- Let us seek CLs using the Direct method, 2nd-order multipliers [C., Oberlack (2014)]:

$$\Lambda_\sigma = \Lambda_\sigma(45 \text{ variables});$$

$$\Lambda_\sigma R^\sigma \equiv \frac{\partial \Phi^i}{\partial x^i} = 0.$$

Euler equations:

$$\nabla \cdot \mathbf{u} = 0,$$

$$\mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u} + \nabla p = 0.$$

Conservation of **generalized momentum**:

- x-direction:

$$\begin{aligned} & \frac{\partial}{\partial t}(f(t)u^1) + \frac{\partial}{\partial x} \left((u^1 f(t) - x f'(t))u^1 + f(t)p \right) \\ & + \frac{\partial}{\partial y} \left((u^1 f(t) - x f'(t))u^2 \right) + \frac{\partial}{\partial z} \left((u^1 f(t) - x f'(t))u^3 \right) = 0. \end{aligned}$$

- Multipliers:

$$\Lambda_1 = f(t)u^1 - x f'(t), \quad \Lambda_2 = f(t), \quad \Lambda_3 = \Lambda_4 = 0.$$

- Arbitrary $f(t)$.
- Similar in y -, z -directions.

Euler equations:

$$\nabla \cdot \mathbf{u} = 0,$$

$$\mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u} + \nabla p = 0.$$

Conservation of angular momentum $\mathbf{x} \times \mathbf{u}$:

- x-direction:

$$\begin{aligned} & \frac{\partial}{\partial t}(zu^2 - yu^3) + \frac{\partial}{\partial x}((zu^2 - yu^3)u^1) \\ & + \frac{\partial}{\partial y}((zu^2 - yu^3)u^2 + zp) + \frac{\partial}{\partial z}((zu^2 - yu^3)u^3 - yp) = 0. \end{aligned}$$

- Multipliers:

$$\Lambda_1 = u_z^2 - u_y^3, \quad \Lambda_2 = 0, \quad \Lambda_3 = z, \quad \Lambda_4 = -y.$$

- Similar in y-, z-directions.

Euler equations:

$$\nabla \cdot \mathbf{u} = 0,$$

$$\mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u} + \nabla p = 0.$$

Conservation of **kinetic energy**:

- x-direction:

$$\frac{\partial}{\partial t} K + \nabla \cdot ((K + p)\mathbf{u}) = 0, \quad K = \frac{1}{2} |\mathbf{u}|^2.$$

- Multipliers:

$$\Lambda_1 = K + p, \quad \Lambda_i = u^i, \quad i = 1, 2, 3.$$

Euler equations:

$$\nabla \cdot \mathbf{u} = 0,$$

$$\mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u} + \nabla p = 0.$$

Generalized continuity equation:

- For arbitrary $k(t)$:

$$\nabla \cdot (k(t) \mathbf{u}) = 0.$$

- Multipliers:

$$\Lambda_1 = k(t), \quad \Lambda_2 = \Lambda_3 = \Lambda_4 = 0.$$

- Arbitrary $k(t)$.

Euler equations:

$$\nabla \cdot \mathbf{u} = 0,$$

$$\mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u} + \nabla p = 0.$$

Conservation of helicity:

- Vorticity: $\boldsymbol{\omega} = \text{curl } \mathbf{u}$.
- Helicity: $h = \mathbf{u} \cdot \boldsymbol{\omega}$.
- Helicity conservation law:

$$\frac{\partial}{\partial t} h + \nabla \cdot (\mathbf{u} \times \nabla E + (\boldsymbol{\omega} \times \mathbf{u}) \times \mathbf{u}) = 0,$$

where $E = K + p$ is the total energy density.

- Topological significance/vortex line linkage.
- Multipliers:

$$\Lambda_1 = 0, \quad \Lambda_i = \omega^i, \quad i = 1, 2, 3.$$

Euler equations:

$$\nabla \cdot \mathbf{u} = 0,$$

$$\mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u} + \nabla p = 0.$$

Vorticity system: conservation of **vorticity**.

- Vorticity: $\boldsymbol{\omega} = \text{curl } \mathbf{u}$.
- Vorticity equations:

$$\text{div } \boldsymbol{\omega} = 0, \quad \boldsymbol{\omega}_t + \text{curl}(\boldsymbol{\omega} \times \mathbf{u}) = 0.$$

Euler equations:

$$\nabla \cdot \mathbf{u} = 0,$$

$$\mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u} + \nabla p = 0.$$

Vorticity system: potential vorticity.

- Vorticity equations:

$$\operatorname{div} \boldsymbol{\omega} = 0, \quad \boldsymbol{\omega}_t + \operatorname{curl}(\boldsymbol{\omega} \times \mathbf{u}) = 0.$$

- CL:

$$(\boldsymbol{\omega} \cdot \nabla F)_t + \nabla \cdot (\boldsymbol{\beta} \times \nabla F - F_t \boldsymbol{\omega}) = 0, \quad \boldsymbol{\beta} \equiv \boldsymbol{\omega} \times \mathbf{u}.$$

- Multipliers:

$$\Lambda_1 = -D_t F, \quad \Lambda_2 = D_x F, \quad \Lambda_3 = D_y F, \quad \Lambda_4 = D_z F,$$

holding for an arbitrary differential function $F = F[\mathbf{u}, p]$.

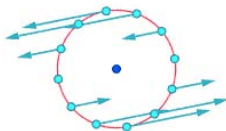
- Details [Müller (1995)], generalizations: [C. & Oberlack (2014)].

Euler classical two-component plane flow:

$$u^z = \omega^x = \omega^y = 0; \quad \frac{\partial}{\partial z} = 0.$$

$$\begin{cases} (u^x)_x + (u^y)_y = 0, \\ (u^x)_t + u^x(u^x)_x + u^y(u^x)_y = -p_x, \\ (u^y)_t + u^x(u^y)_x + u^y(u^y)_y = -p_y; \end{cases}$$

$$\begin{cases} \omega^z + (u^x)_y - (u^y)_x = 0, \\ (\omega^z)_t + u^x(\omega^z)_x + u^y(\omega^z)_y = 0. \end{cases}$$



Euler classical two-component plane flow:

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Enstrophy Conservation

- Enstrophy: $\mathcal{E} = |\omega|^2 = (\omega^z)^2$.
- Material conservation law:

$$\frac{d}{dt} \mathcal{E} = D_t \mathcal{E} + D_x (u^x \mathcal{E}) + D_y (u^y \mathcal{E}) = 0.$$

- Was only known to hold for plane flows, $(2 + 1)$ -dimensions.

Euler classical two-component plane flow:

$$u^z = \omega^x = \omega^y = 0; \quad \frac{\partial}{\partial z} = 0.$$

$$\begin{cases} (u^x)_x + (u^y)_y = 0, \\ (u^x)_t + u^x(u^x)_x + u^y(u^x)_y = -p_x, \\ (u^y)_t + u^x(u^y)_x + u^y(u^y)_y = -p_y; \end{cases}$$

$$\begin{cases} \omega^z + (u^x)_y - (u^y)_x = 0, \\ (\omega^z)_t + u^x(\omega^z)_x + u^y(\omega^z)_y = 0. \end{cases}$$

Other Plane Flow CLs

- Several additional vorticity-related CLs known for plane flows (e.g., [Batchelor (2000)]);

Navier-Stokes Equations equations in 3 + 1 dimensions

$$\begin{aligned}\nabla \cdot \mathbf{u} &= 0, \\ \mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u} + \nabla p - \nu \nabla^2 \mathbf{u} &= 0.\end{aligned}$$

Vorticity formulation:

$$\begin{aligned}\nabla \cdot \mathbf{u} &= 0, \quad \boldsymbol{\omega} = \nabla \times \mathbf{u}, \\ \boldsymbol{\omega}_t + \nabla \times (\boldsymbol{\omega} \times \mathbf{u}) - \nu \nabla^2 \boldsymbol{\omega} &= 0.\end{aligned}$$

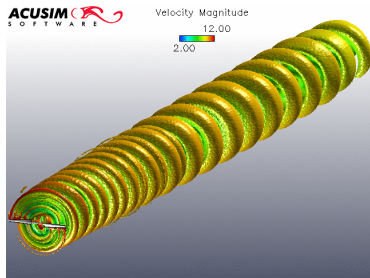
Basic conservation laws:

- Momentum / generalized momentum: $\Theta = f(t)u^i, \quad i = 1, 2, 3.$
- Angular momentum: $\Theta = (\mathbf{r} \times \mathbf{u})^i, \quad i = 1, 2, 3.$
- Vorticity: $\Theta = \omega^i, \quad i = 1, 2, 3.$
- Potential vorticity.

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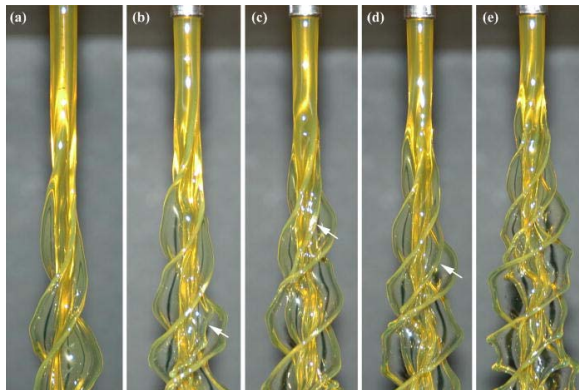
Examples of Helical Flows in Nature

- Wind turbine wakes in aerodynamics [*Vermeer, Sorensen & Crespo, 2003*]



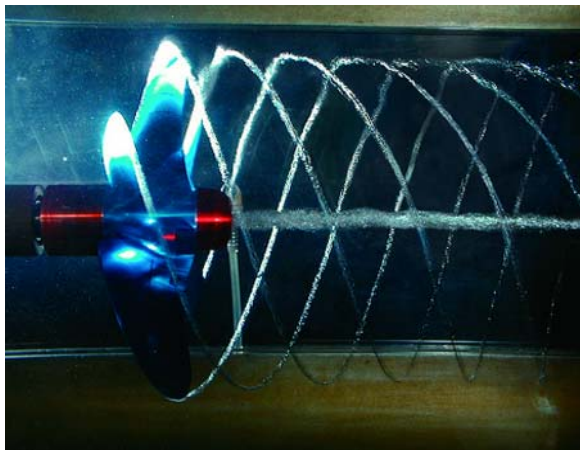
Examples of Helical Flows in Nature

- Helical instability of rotating viscous jets [*Kubitschek & Weidman, 2007*]



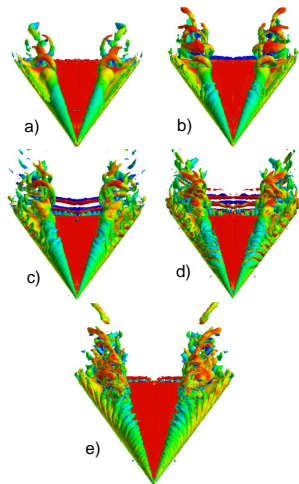
Examples of Helical Flows in Nature

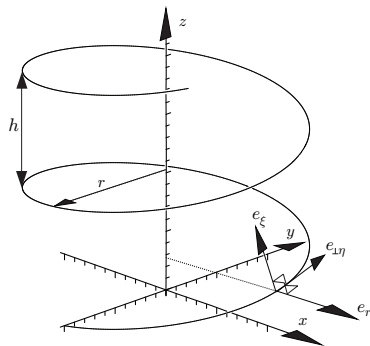
- Helical water flow past a propeller



Examples of Helical Flows in Nature

- Wing tip vortices, in particular, on delta wings [Mitchell, Morton & Forsythe, 1997]

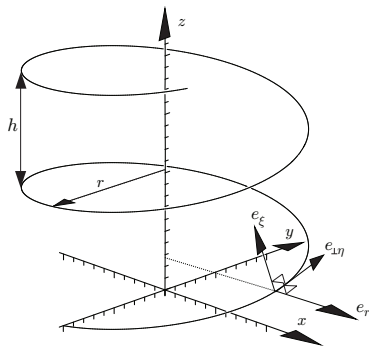




Helical Coordinates

- Cylindrical coordinates: (r, φ, z) . **Helical coordinates: (r, η, ξ)**

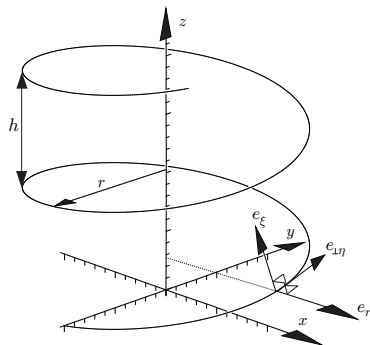
$$\xi = az + b\varphi, \quad \eta = a\varphi - b\frac{z}{r^2}, \quad a, b = \text{const}, \quad a^2 + b^2 > 0.$$



Orthogonal Basis

$$\mathbf{e}_r = \frac{\nabla r}{|\nabla r|}, \quad \mathbf{e}_\xi = \frac{\nabla \xi}{|\nabla \xi|}, \quad \mathbf{e}_{\perp\eta} = \frac{\nabla_{\perp\eta}}{|\nabla_{\perp\eta}|} = \mathbf{e}_\xi \times \mathbf{e}_r.$$

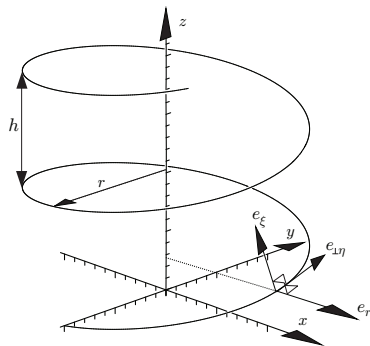
- Scaling factors: $H_r = 1, H_\eta = r, H_\xi = B(r), \quad B(r) = \frac{r}{\sqrt{a^2 r^2 + b^2}}.$



Vector expansion

$$\mathbf{u} = u^r \mathbf{e}_r + u^\varphi \mathbf{e}_\varphi + u^z \mathbf{e}_z = u^r \mathbf{e}_r + u^\eta \mathbf{e}_{\perp\eta} + u^\xi \mathbf{e}_\xi.$$

$$u^\eta = \mathbf{u} \cdot \mathbf{e}_{\perp\eta} = B \left(a u^\varphi - \frac{b}{r} u^z \right), \quad u^\xi = \mathbf{u} \cdot \mathbf{e}_\xi = B \left(\frac{b}{r} u^\varphi + a u^z \right).$$



Helical invariance: generalizes axial and translational invariance

- Helical coordinates: r , $\xi = az + b\varphi$, $\eta = a\varphi - bz/r^2$.
- **General helical symmetry:** $f = f(r, \xi)$, $a, b \neq 0$.
- **Axial:** $a = 1$, $b = 0$. **z-Translational:** $a = 0$, $b = 1$.

Navier-Stokes Equations:

$$\begin{aligned}\nabla \cdot \mathbf{u} &= 0, \\ \mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u} + \nabla p - \nu \nabla^2 \mathbf{u} &= 0.\end{aligned}$$

Navier-Stokes Equations:

$$\begin{aligned}\nabla \cdot \mathbf{u} &= 0, \\ \mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u} + \nabla p - \nu \nabla^2 \mathbf{u} &= 0.\end{aligned}$$

Continuity:

$$\frac{1}{r} u^r + (u^r)_r + \frac{1}{B} (u^\xi)_\xi = 0$$

Navier-Stokes Equations:

$$\begin{aligned}\nabla \cdot \mathbf{u} &= 0, \\ \mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u} + \nabla p - \nu \nabla^2 \mathbf{u} &= 0.\end{aligned}$$

r -momentum:

$$\begin{aligned}(u^r)_t + u^r(u^r)_r + \frac{1}{B} u^\xi (u^r)_\xi - \frac{B^2}{r} \left(\frac{b}{r} u^\xi + a u^\eta \right)^2 &= -p_r \\ + \nu \left[\frac{1}{r} (r(u^r)_r)_r + \frac{1}{B^2} (u^r)_{\xi\xi} - \frac{1}{r^2} u^r - \frac{2bB}{r^2} \left(a(u^\eta)_\xi + \frac{b}{r} (u^\xi)_\xi \right) \right] &\end{aligned}$$

Navier-Stokes Equations:

$$\begin{aligned}\nabla \cdot \mathbf{u} &= 0, \\ \mathbf{u}_t + (\mathbf{u} \cdot \nabla)\mathbf{u} + \nabla p - \nu \nabla^2 \mathbf{u} &= 0.\end{aligned}$$

η -momentum:

$$\begin{aligned}(u^\eta)_t + u^r (u^\eta)_r + \frac{1}{B} u^\xi (u^\eta)_\xi + \frac{a^2 B^2}{r} u^r u^\eta \\ = \nu \left[\frac{1}{r} (r(u^\eta)_r)_r + \frac{1}{B^2} (u^\eta)_{\xi\xi} + \frac{a^2 B^2 (a^2 B^2 - 2)}{r^2} u^\eta + \frac{2abB}{r^2} \left((u^r)_\xi - (Bu^\xi)_r \right) \right]\end{aligned}$$

Navier-Stokes Equations:

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ξ -momentum:

$$\begin{aligned}(u^\xi)_t + u^r (u^\xi)_r + \frac{1}{B} u^\xi (u^\xi)_\xi + \frac{2abB^2}{r^2} u^r u^\eta + \frac{b^2 B^2}{r^3} u^r u^\xi &= -\frac{1}{B} p_\xi \\ + \nu \left[\frac{1}{r} (r(u^\xi)_r)_r + \frac{1}{B^2} (u^\xi)_{\xi\xi} + \frac{a^4 B^4 - 1}{r^2} u^\xi + \frac{2bB}{r} \left(\frac{b}{r^2} (u^r)_\xi + \left(\frac{aB}{r} u^\eta \right)_r \right) \right] &\end{aligned}$$

Navier-Stokes Equations, Vorticity Formulation:

$$\nabla \cdot \mathbf{u} = 0,$$

$$\nabla \times \mathbf{u} =: \boldsymbol{\omega} = \omega^r \mathbf{e}_r + \omega^\eta \mathbf{e}_{\perp\eta} + \omega^\xi \mathbf{e}_\xi,$$

$$\boldsymbol{\omega}_t + \nabla \times (\boldsymbol{\omega} \times \mathbf{u}) - \nu \nabla^2 \boldsymbol{\omega} = 0.$$

Navier-Stokes Equations, Vorticity Formulation:

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Vorticity definition:

$$\begin{aligned}\omega^r &= -\frac{1}{B}(u^\eta)_\xi, \\ \omega^\eta &= \frac{1}{B}(u^r)_\xi - \frac{1}{r}(ru^\xi)_r - \frac{2abB^2}{r^2}u^\eta + \frac{a^2B^2}{r}u^\xi, \\ \omega^\xi &= (u^\eta)_r + \frac{a^2B^2}{r}u^\eta\end{aligned}$$

Navier-Stokes Equations, Vorticity Formulation:

$$\begin{aligned}\nabla \cdot \mathbf{u} &= 0, \\ \nabla \times \mathbf{u} &=: \boldsymbol{\omega} = \omega^r \mathbf{e}_r + \omega^\eta \mathbf{e}_{\perp\eta} + \omega^\xi \mathbf{e}_\xi, \\ \boldsymbol{\omega}_t + \nabla \times (\boldsymbol{\omega} \times \mathbf{u}) - \nu \nabla^2 \boldsymbol{\omega} &= 0.\end{aligned}$$

r -Momentum:

$$\begin{aligned}(\omega^r)_t + u_r (\omega^r)_r + \frac{1}{B} u^\xi (\omega^r)_\xi &= \omega^r (u^r)_r + \frac{1}{B} \omega^\xi (u^r)_\xi \\ + \nu \left[\frac{1}{r} (r(\omega^r)_r)_r + \frac{1}{B^2} (\omega^r)_{\xi\xi} - \frac{1}{r^2} \omega^r - \frac{2bB}{r^2} \left(a(\omega^\eta)_\xi + \frac{b}{r} (\omega^\xi)_\xi \right) \right]\end{aligned}$$

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$$\begin{aligned}\nabla \cdot \mathbf{u} &= 0, \\ \nabla \times \mathbf{u} &=: \boldsymbol{\omega} = \omega^r \mathbf{e}_r + \omega^\eta \mathbf{e}_{\perp\eta} + \omega^\xi \mathbf{e}_\xi, \\ \boldsymbol{\omega}_t + \nabla \times (\boldsymbol{\omega} \times \mathbf{u}) - \nu \nabla^2 \boldsymbol{\omega} &= 0.\end{aligned}$$

η -Momentum:

$$\begin{aligned}(\omega^\eta)_t + u^r (\omega^\eta)_r + \frac{1}{B} u^\xi (\omega^\eta)_\xi \\ - \frac{a^2 B^2}{r} (u^r \omega^\eta - u^\eta \omega^r) + \frac{2abB^2}{r^2} (u^\xi \omega^r - u^r \omega^\xi) = \omega^r (u^\eta)_r + \frac{1}{B} \omega^\xi (u^\eta)_\xi \\ + \nu \left[\frac{1}{r} (r(\omega^\eta)_r)_r + \frac{1}{B^2} (\omega^\eta)_{\xi\xi} + \frac{a^2 B^2 (a^2 B^2 - 2)}{r^2} \omega^\eta + \frac{2abB}{r^2} \left((\omega^r)_\xi - (B\omega^\xi)_r \right) \right]\end{aligned}$$

Navier-Stokes Equations, Vorticity Formulation:

$$\begin{aligned}\nabla \cdot \mathbf{u} &= 0, \\ \nabla \times \mathbf{u} &=: \boldsymbol{\omega} = \omega^r \mathbf{e}_r + \omega^\eta \mathbf{e}_{\perp\eta} + \omega^\xi \mathbf{e}_\xi, \\ \boldsymbol{\omega}_t + \nabla \times (\boldsymbol{\omega} \times \mathbf{u}) - \nu \nabla^2 \boldsymbol{\omega} &= 0.\end{aligned}$$

ξ -Momentum:

$$\begin{aligned}(\omega^\xi)_t + u^r (\omega^\xi)_r + \frac{1}{B} u^\xi (\omega^\xi)_\xi \\ + \frac{1 - a^2 B^2}{r} (u^\xi \omega^r - u^r \omega^\xi) = \omega^r (u^\xi)_r + \frac{1}{B} \omega^\xi (u^\xi)_\xi \\ + \nu \left[\frac{1}{r} (r (\omega^\xi)_r)_r + \frac{1}{B^2} (\omega^\xi)_{\xi\xi} + \frac{a^4 B^4 - 1}{r^2} \omega^\xi + \frac{2bB}{r} \left(\frac{b}{r^2} (\omega^r)_\xi + \left(\frac{aB}{r} \omega^\eta \right)_r \right) \right]\end{aligned}$$

For helically symmetric flows:

- Seek local conservation laws

$$\frac{\partial \Theta}{\partial t} + \nabla \cdot \Phi \equiv \frac{\partial \Theta}{\partial t} + \frac{1}{r} \frac{\partial}{\partial r} (r \Phi^r) + \frac{1}{B} \frac{\partial \Phi^\xi}{\partial \xi} = 0$$

using divergence expressions

$$\frac{\partial \Gamma^1}{\partial t} + \frac{\partial \Gamma^2}{\partial r} + \frac{\partial \Gamma^3}{\partial \xi} = r \left[\frac{\partial}{\partial t} \left(\frac{\Gamma^1}{r} \right) + \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\Gamma^2}{r} \right) + \frac{1}{B} \frac{\partial}{\partial \xi} \left(\frac{B}{r} \Gamma^3 \right) \right] = 0,$$

i.e.,

$$\Theta \equiv \frac{\Gamma^1}{r}, \quad \Phi^r \equiv \frac{\Gamma^2}{r}, \quad \Phi^\xi \equiv \frac{B}{r} \Gamma^3.$$

- 1st-order multipliers in primitive variables.
- 0th-order multipliers in vorticity formulation.

Primitive variables - EP1 - Kinetic energy

$$\Theta = K, \quad \Phi^r = u^r(K + p), \quad \Phi^\xi = u^\xi(K + p), \quad K = \frac{1}{2}|\mathbf{u}|^2.$$

Primitive variables - EP2 - z-momentum

$$\Theta = B \left(-\frac{b}{r}u^\eta + au^\xi \right) = u^z, \quad \Phi^r = u^r u^z, \quad \Phi^\xi = u^\xi u^z + aBp.$$

Primitive variables - EP3 - z-angular momentum

$$\Theta = rB \left(au^\eta + \frac{b}{r}u^\xi \right) = ru^\varphi, \quad \Phi^r = ru^r u^\varphi, \quad \Phi^\xi = ru^\xi u^\varphi + bBp.$$

Primitive variables - EP4 - Generalized momenta/angular momenta

$$\Theta = F \left(\frac{r}{B}u^\eta \right), \quad \Phi^r = u^r F \left(\frac{r}{B}u^\eta \right), \quad \Phi^\xi = u^\xi F \left(\frac{r}{B}u^\eta \right),$$

where $F(\cdot)$ is an arbitrary function.

Vorticity formulation - EV1 - Conservation of helicity

Helicity:

$$h = \mathbf{u} \cdot \boldsymbol{\omega} = u^r \omega^r + u^\eta \omega^\eta + u^\xi \omega^\xi.$$

The conservation law:

$$\Theta = h,$$

$$\Phi^r = \omega^r \left(E - (u^\eta)^2 - (u^\xi)^2 \right) + u^r (h - u^r \omega^r),$$

$$\Phi^\xi = \omega^\xi \left(E - (u^r)^2 - (u^\eta)^2 \right) + u^\xi (h - u^\xi \omega^\xi),$$

where

$$E = \frac{1}{2} |\mathbf{u}|^2 + p = \frac{1}{2} \left((u^r)^2 + (u^\eta)^2 + (u^\xi)^2 \right) + p$$

is the total energy density. In vector notation:

$$\frac{\partial}{\partial t} h + \nabla \cdot (\mathbf{u} \times \nabla E + (\boldsymbol{\omega} \times \mathbf{u}) \times \mathbf{u}) = 0.$$

Vorticity formulation - EV2 - Generalized helicity

Helicity:

$$h = \mathbf{u} \cdot \boldsymbol{\omega} = u^r \omega^r + u^\eta \omega^\eta + u^\xi \omega^\xi.$$

$$\frac{\partial}{\partial t} \left(h H \left(\frac{r}{B} u^\eta \right) \right) + \nabla \cdot \left[H \left(\frac{r}{B} u^\eta \right) [\mathbf{u} \times \nabla E + (\boldsymbol{\omega} \times \mathbf{u}) \times \mathbf{u}] + E u^\eta \mathbf{e}_{\perp \eta} \times \nabla H \left(\frac{r}{B} u^\eta \right) \right] = 0$$

for an arbitrary function $H = H(\cdot)$.

Vorticity formulation - EV3 - Vorticity conservation laws

$$\Theta = \frac{Q(t)}{r} \omega^\varphi,$$

$$\Phi^r = \frac{1}{r} (Q(t)[u^r \omega^\varphi - \omega^r u^\varphi] + Q'(t)u^z),$$

$$\Phi^\xi = -\frac{aB}{r} (Q(t)[u^\eta \omega^\xi - u^\xi \omega^\eta] + Q'(t)u^r),$$

where $Q(t)$ is an arbitrary function.

Vorticity formulation - EV4 - Vorticity conservation law

$$\Theta = -rB \left(a^3 \omega^\eta - \frac{b^3}{r^3} \omega^\xi \right),$$

$$\Phi^r = -2a^2 u^r u^z - a^3 B r (u^r \omega^\eta - u^\eta \omega^r) + \frac{B b^3}{r^2} (u^r \omega^\xi - u^\xi \omega^r),$$

$$\Phi^\xi = a^3 B [(u^r)^2 + (u^\eta)^2 - (u^\xi)^2 + r(u^\eta \omega^\xi - u^\xi \omega^\eta)] + \frac{2a^2 b B}{r} u^\eta u^\xi.$$

Vorticity formulation - EV5 - Vorticity conservation law

$$\Theta = -\frac{B}{r^2} \left(\frac{b^2 r^2}{B^2} \omega^\xi + a^3 r^4 \left(-\frac{b}{r} \omega^\eta + a \omega^\xi \right) \right) = -\frac{B}{r^2} \left(\frac{b^2 r^2}{B^2} \omega^\xi + \frac{a^3 r^4}{B} \omega^z \right),$$

$$\Phi^r = a^3 r B \left(2u^r \left(a u^\eta + \frac{b}{r} u^\xi \right) + b(u^r \omega^\eta - u^\eta \omega^r) \right) \\ - \frac{a^4 r^4 + a^2 r^2 b^2 + b^4}{r \sqrt{a^2 r^2 + b^2}} (u^r \omega^\xi - u^\xi \omega^r),$$

$$\Phi^\xi = -a^3 b B \left((u^r)^2 + (u^\eta)^2 - (u^\xi)^2 + r(u^\eta \omega^\xi - u^\xi \omega^\eta) \right) + 2a^4 r B u^\eta u^\xi.$$

Vorticity formulation - EV6 - Vorticity conservation law

$$\nabla \cdot \Phi = 0, \quad \Phi^r = N \omega^r - \frac{1}{B} N_\xi u^\eta, \quad \Phi^\xi = N \omega^\xi,$$

for an arbitrary $N(t, \xi)$.

- Generalization of the obvious divergence expression $\nabla \cdot (G(t)\omega) = 0$.

Primitive variables - NSP1 - z-momentum.

$$\Theta = u^z, \quad \Phi^r = u^r u^z - \nu (u^z)_r, \quad \Phi^\xi = u^\xi u^z + aBp - \frac{\nu}{B} (u^z)_\xi.$$

Primitive variables - NSP2 - generalized momentum

$$\Theta = \frac{r}{B} u^\eta,$$

$$\begin{aligned} \Phi^r &= \frac{r}{B} u^r u^\eta - \nu \left[-2aB \left(au^\eta + 2\frac{b}{r} u^\xi \right) + \left(\frac{r}{B} u^\eta \right)_r \right] \\ &= \frac{r}{B} u^r u^\eta - \nu \left[-2au^\varphi + \left(\frac{r}{B} u^\eta \right)_r \right], \end{aligned}$$

$$\Phi^\xi = \frac{r}{B} u^\eta u^\xi - \nu \frac{1}{B} \left[\frac{2abB^2}{r} u^r + \left(\frac{r}{B} u^\eta \right)_\xi \right].$$

Vorticity formulation - NSV1 - Family of vorticity conservation laws

$$\Theta = \frac{Q(t)}{r} B \left(a\omega^\eta + \frac{b}{r}\omega^\xi \right) = \frac{Q(t)}{r} \omega^\varphi,$$

$$\Phi^r = \frac{1}{r} \left\{ Q(t) \left[u^r B \left(a\omega^\eta + \frac{b}{r}\omega^\xi \right) - \omega^r B \left(au^\eta + \frac{b}{r}u^\xi \right) \right] + Q'(t) B \left(-\frac{b}{r}u^\eta + au^\xi \right) - Q(t)\nu \left[\frac{aB}{r}\omega^\eta + \frac{b^2B}{r(a^2r^2 + b^2)} \left(a\omega^\eta + \frac{b}{r}\omega^\xi \right) + B \left(a\omega_r^\eta + \frac{b}{r}\omega_r^\xi \right) \right] \right\},$$

$$\Phi^\xi = -\frac{B}{r} \left\{ aQ(t) [u^\eta\omega^\xi - u^\xi\omega^\eta] + aQ'(t)u^r + \frac{Q(t)}{r^3} \nu \left[\frac{r^3}{B} \left(a\omega_\xi^\eta + \frac{b}{r}\omega_\xi^\xi \right) + 2br\omega^r \right] \right\},$$

for an arbitrary function $Q(t)$.

Vorticity formulation - NSV2 - Vorticity conservation law

$$\Theta = -rB \left(a^3 \omega^\eta - \frac{b^3}{r^3} \omega^\xi \right),$$

$$\begin{aligned} \Phi^r = & -\frac{B}{r^2} \left(a^3 r^3 (u^r \omega^\eta - u^\eta \omega^r) - b^3 (u^r \omega^\xi - u^\xi \omega^r) \right) - 2a^2 B u^r \left(-\frac{b}{r} u^\eta + a u^\xi \right) \\ & - \frac{B}{r^2} \nu \left[\frac{r^2}{B^2} \left(a \omega^\eta + \frac{b}{r} \omega^\xi \right) - r^3 \left(a^3 \omega_r^\eta - \frac{b^3}{r^3} \omega_r^\xi \right) + abB^2 r \left(\frac{b^3}{r^3} \omega^\eta + a^3 \omega^\xi \right) \right], \end{aligned}$$

$$\begin{aligned} \Phi^\xi = & a^3 B \left((u^r)^2 + (u^\eta)^2 - (u^\xi)^2 + r (u^\eta \omega^\xi - u^\xi \omega^\eta) \right) + \frac{2a^2 b B}{r} u^\eta u^\xi \\ & + \frac{2a^2 b B}{r} \nu \left[\left(1 - \frac{b^2}{a^2 r^2} \right) \omega^r + \frac{r^2}{2a^2 b B} \left(a^3 \omega_\xi^\eta - \frac{b^3}{r^3} \omega_\xi^\xi \right) \right]. \end{aligned}$$

Vorticity formulation - NSV3 - Vorticity conservation law

$$\Theta = -\frac{B}{r^2} \left(\frac{b^2 r^2}{B^2} \omega^\xi + a^3 r^4 \left(-\frac{b}{r} \omega^\eta + a \omega^\xi \right) \right) = -\frac{B}{r^2} \left(\frac{b^2 r^2}{B^2} \omega^\xi + \frac{a^3 r^4}{B} \omega^z \right),$$

$$\begin{aligned} \Phi^r = & a^3 r B \left(2u^r \left(au^\eta + \frac{b}{r} u^\xi \right) + b(u^r \omega^\eta - u^\eta \omega^r) \right) \\ & - \frac{a^4 r^4 + a^2 r^2 b^2 + b^4}{r \sqrt{a^2 r^2 + b^2}} (u^r \omega^\xi - u^\xi \omega^r) \\ & + \nu \left[4a^3 B \left(au^\eta + \frac{b}{r} u^\xi \right) - a^3 b r B (\omega^\eta)_r + \frac{B}{r^3} \left(b^4 - a^4 r^4 - \frac{a^6 r^6}{a^2 r^2 + b^2} \right) \omega^\xi \right. \\ & \left. + \frac{B}{r^2} (a^4 r^4 + a^2 r^2 b^2 + b^4) (\omega^\xi)_r + \frac{ab}{B} \left(2 + \frac{a^4 r^4}{(a^2 r^2 + b^2)^2} \right) \omega^\eta \right], \end{aligned}$$

$$\begin{aligned} \Phi^\xi = & -a^3 b B \left((u^r)^2 + (u^\eta)^2 - (u^\xi)^2 + r(u^\eta \omega^\xi - u^\xi \omega^\eta) \right) + 2a^4 r B u^\eta u^\xi \\ & + \nu \left[\frac{1}{r^2} (a^4 r^4 + a^2 r^2 b^2 + b^4) (\omega^\xi)_\xi - a^3 b r (\omega^\eta)_\xi - \frac{4a^3 b B}{r} u^r + \frac{2b^4 B}{r^3} \omega^r \right]. \end{aligned}$$

Generalized enstrophy for inviscid plane flow (known)

$$\Theta = N(\omega^z), \quad \Phi^x = u^x N(\omega^z), \quad \Phi^y = u^y N(\omega^z),$$

for an arbitrary $N(\cdot)$, equivalent to a material conservation law

$$\frac{d}{dt} N(\omega^z) = 0.$$

Generalized enstrophy for inviscid plane flow (known)

$$\Theta = N(\omega^z), \quad \Phi^x = u^x N(\omega^z), \quad \Phi^y = u^y N(\omega^z),$$

for an arbitrary $N(\cdot)$, equivalent to a material conservation law

$$\frac{d}{dt} N(\omega^z) = 0.$$

Generalized enstrophy for inviscid axisymmetric flow

$$\Theta = S\left(\frac{1}{r}\omega^\varphi\right), \quad \Phi^r = u^r S\left(\frac{1}{r}\omega^\varphi\right), \quad \Phi^z = u^z S\left(\frac{1}{r}\omega^\varphi\right)$$

for arbitrary $S(\cdot)$.

Generalized enstrophy for inviscid plane flow (known)

$$\Theta = N(\omega^z), \quad \Phi^x = u^x N(\omega^z), \quad \Phi^y = u^y N(\omega^z),$$

for an arbitrary $N(\cdot)$, equivalent to a material conservation law

$$\frac{d}{dt} N(\omega^z) = 0.$$

Generalized enstrophy for inviscid axisymmetric flow

$$\Theta = S\left(\frac{1}{r}\omega^\varphi\right), \quad \Phi^r = u^r S\left(\frac{1}{r}\omega^\varphi\right), \quad \Phi^z = u^z S\left(\frac{1}{r}\omega^\varphi\right)$$

for arbitrary $S(\cdot)$.

- Several additional conservation laws arise for **plane** and **axisymmetric**, **inviscid** and **viscous** flows (details in paper).

Helically-Invariant Equations

- Full three-component Euler and Navier-Stokes equations written in helically-invariant form.
- Two-component reductions.

Additional Conservation Laws

- Three-component Euler:
 - Generalized momenta. Generalized helicity. Additional vorticity CLs.
- Three-component Navier-Stokes:
 - Additional CLs in primitive and vorticity formulation.
- Two-component flows:
 - Infinite set of enstrophy-related vorticity CLs (inviscid case).
 - Additional CLs in viscous and inviscid case, for plane and axisymmetric flows.

- 1 Fluid Dynamics Equations
- 2 CLs of Constant-Density Euler and N-S Equations
- 3 CLs of Helically Invariant Flows
- 4 CLs of An Inviscid Model in Gas Dynamics**
- 5 CLs of a Surfactant Flow Model
- 6 Discussion

- Euler equations:

$$\rho_t + \nabla \cdot (\rho \mathbf{u}) = 0, \quad \rho(\mathbf{u}_t + (\mathbf{u} \cdot \nabla) \mathbf{u}) + \nabla p = 0.$$

- A CL classification for 2D, 3D barotropic model:

$$p = p(\rho) \quad (S = \text{const}).$$

[Anco & Dar (2010)].

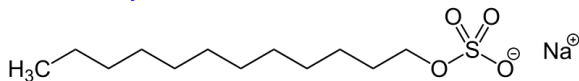
- 1 Fluid Dynamics Equations
- 2 CLs of Constant-Density Euler and N-S Equations
- 3 CLs of Helically Invariant Flows
- 4 CLs of An Inviscid Model in Gas Dynamics
- 5 CLs of a Surfactant Flow Model**
- 6 Discussion

Surfactants

- "Surfactant" = "Surface active agent".
- Act as detergents, wetting agents, emulsifiers, foaming agents, and dispersants.
- Consist of a hydrophobic group (*tail*) and a hydrophilic group (*head*).

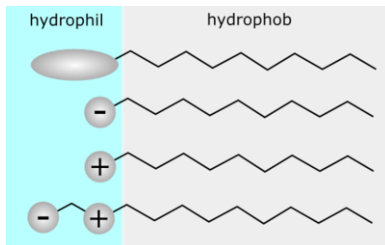
Surfactants

- "Surfactant" = "Surface active agent".
- Act as detergents, wetting agents, emulsifiers, foaming agents, and dispersants.
- Consist of a hydrophobic group (*tail*) and a hydrophilic group (*head*).
- Sodium lauryl sulfate:

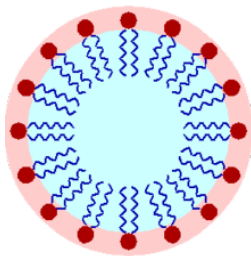


Surfactants

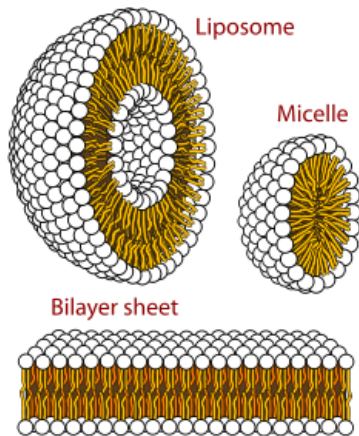
- "Surfactant" = "Surface active agent".
- Act as detergents, wetting agents, emulsifiers, foaming agents, and dispersants.
- Consist of a hydrophobic group (*tail*) and a hydrophilic group (*head*).
- Hydrophilic groups (heads) can have various properties:



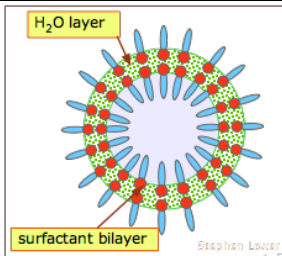
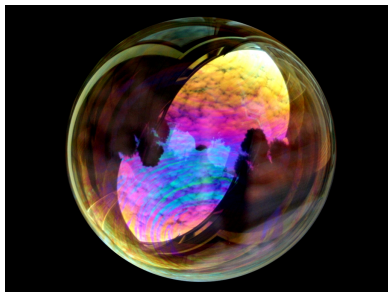
- Surfactant molecules adsorb at phase separation interfaces.
 - Stabilization of growth of bubbles / droplets.
 - Creation of emulsions of insoluble substances.
 - Multiple industrial and medical applications.



- Can form micelles, double layers, etc.



- Soap bubbles...

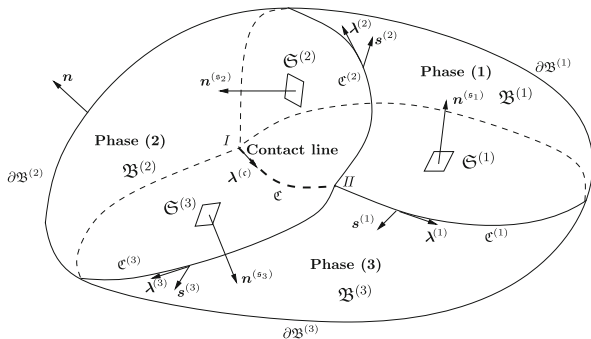


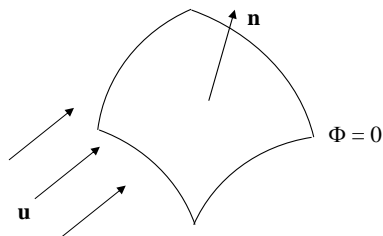
Derivation:

Can be derived as a special case of multiphase flows with moving interfaces and contact lines:

[Y.Wang, M. Oberlack, 2011]

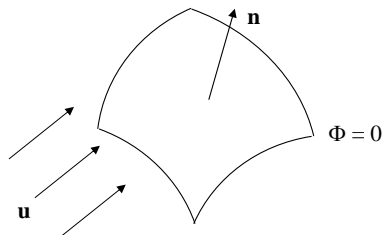
- Illustration:





Parameters

- Surfactant concentration $c = c(\mathbf{x}, t)$.
- Flow velocity $\mathbf{u}(\mathbf{x}, t)$.
- Two-phase interface: phase separation surface $\Phi(\mathbf{x}, t) = 0$.
- Unit normal: $\mathbf{n} = -\frac{\nabla\Phi}{|\nabla\Phi|}$.

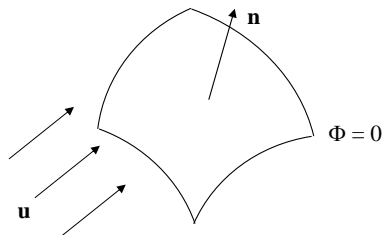


Surface gradient

- Surface projection tensor: $p_{ij} = \delta_{ij} - n_i n_j$.
- Surface gradient operator: $\nabla^s = \mathbf{p} \cdot \nabla = (\delta_{ij} - n_i n_j) \frac{\partial}{\partial x^j}$.
- Surface Laplacian:

$$\Delta^s F = (\delta_{ij} - n_i n_j) \frac{\partial}{\partial x^j} \left((\delta_{ik} - n_i n_k) \frac{\partial F}{\partial x^k} \right).$$

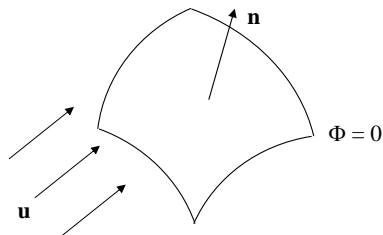
Surfactant Transport Equations (ctd.)



Governing equations

- Incompressibility condition: $\nabla \cdot \mathbf{u} = 0$.
- Fluid dynamics equations: Euler or Navier-Stokes.
- Interface transport by the flow: $\Phi_t + \mathbf{u} \cdot \nabla \Phi = 0$.
- Surfactant transport equation:

$$c_t + u^i \frac{\partial c}{\partial x^i} - cn_i n_j \frac{\partial u^j}{\partial x^i} - \alpha (\delta_{ij} - n_i n_j) \frac{\partial}{\partial x^j} \left((\delta_{ik} - n_i n_k) \frac{\partial c}{\partial x^k} \right) = 0.$$



Fully conserved form

- Specific numerical methods (e.g., discontinuous Galerkin) require the system to be written in a **fully conserved form**.
- Straightforward for continuity, momentum, and interface transport equations.
- **Can the surfactant transport equation be written in the conserved form?**

$$c_t + u^i \frac{\partial c}{\partial x^i} - cn_i n_j \frac{\partial u^i}{\partial x^j} - \alpha (\delta_{ij} - n_i n_j) \frac{\partial}{\partial x^j} \left((\delta_{ik} - n_i n_k) \frac{\partial c}{\partial x^k} \right) = 0.$$

Governing equations ($\alpha = 0$)

$$R^1 = \frac{\partial u^i}{\partial x^i} = 0,$$

$$R^2 = \Phi_t + \frac{\partial(u^i \Phi)}{\partial x^i} = 0,$$

$$R^3 = c_t + u^i \frac{\partial c}{\partial x^i} - cn_i n_j \frac{\partial u^i}{\partial x^j} = 0.$$

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Multiplier ansatz

$$\Lambda^i = \Lambda^i(t, \mathbf{x}, \Phi, c, \mathbf{u}, \partial\Phi, \partial c, \partial\mathbf{u}, \partial^2\Phi, \partial^2 c, \partial^2\mathbf{u}).$$

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Multiplier ansatz

$$\Lambda^i = \Lambda^i(t, \mathbf{x}, \Phi, c, \mathbf{u}, \partial\Phi, \partial c, \partial\mathbf{u}, \partial^2\Phi, \partial^2 c, \partial^2\mathbf{u}).$$

Conservation Law Determining Equations

$$E_{u^j}(\Lambda^\sigma R^\sigma) = 0, \quad j = 1, \dots, 3; \quad E_\Phi(\Lambda^\sigma R^\sigma) = 0; \quad E_c(\Lambda^\sigma R^\sigma) = 0.$$

Governing equations ($\alpha = 0$)

$$R^1 = \frac{\partial u^i}{\partial x^i} = 0,$$

$$R^2 = \Phi_t + \frac{\partial(u^i \Phi)}{\partial x^i} = 0,$$

$$R^3 = c_t + u^i \frac{\partial c}{\partial x^i} - cn_i n_j \frac{\partial u^i}{\partial x^j} = 0.$$

Principal Result 1 (multipliers)

- There exist an **infinite family** of multiplier sets with $\Lambda^3 \neq 0$, i.e., essentially involving c .
- Family of conservation laws with

$$\Lambda^3 = |\nabla \Phi| \mathcal{K}(\Phi, c|\nabla \Phi|).$$

Governing equations ($\alpha = 0$)

$$R^1 = \frac{\partial u^i}{\partial x^i} = 0,$$

$$R^2 = \Phi_t + \frac{\partial(u^i \Phi)}{\partial x^i} = 0,$$

$$R^3 = c_t + u^i \frac{\partial c}{\partial x^i} - cn_i n_j \frac{\partial u^i}{\partial x^j} = 0.$$

Principal Result 1 (divergence expressions)

- Usual form:

$$\frac{\partial}{\partial t} \mathcal{G}(\Phi, c | \nabla \Phi |) + \frac{\partial}{\partial x^i} \left(u^i \mathcal{G}(\Phi, c | \nabla \Phi |) \right) = 0.$$

- Material form:

$$\frac{d}{dt} \mathcal{G}(\Phi, c | \nabla \Phi |) = 0.$$

Governing equations ($\alpha = 0$)

$$R^1 = \frac{\partial u^i}{\partial x^i} = 0,$$

$$R^2 = \Phi_t + \frac{\partial(u^i \Phi)}{\partial x^i} = 0,$$

$$R^3 = c_t + u^i \frac{\partial c}{\partial x^i} - cn_i n_j \frac{\partial u^i}{\partial x^j} = 0.$$

Simplest conservation law with c -dependence

- Can take $\mathcal{G}(\Phi, c|\nabla\Phi|) = c|\nabla\Phi|$.

$$\frac{\partial}{\partial t}(c|\nabla\Phi|) + \frac{\partial}{\partial x^i}(u^i c|\nabla\Phi|) = 0.$$

Governing equations ($\alpha \neq 0$)

$$R^1 = \frac{\partial u^i}{\partial x^i} = 0,$$

$$R^2 = \Phi_t + \frac{\partial(u^i \Phi)}{\partial x^i} = 0,$$

$$R^3 = c_t + u^i \frac{\partial c}{\partial x^i} - cn_i n_j \frac{\partial u^i}{\partial x^j} - \alpha(\delta_{ij} - n_i n_j) \frac{\partial}{\partial x^j} \left((\delta_{ik} - n_i n_k) \frac{\partial c}{\partial x^k} \right) = 0.$$

Governing equations ($\alpha = 0$)

$$R^1 = \frac{\partial u^i}{\partial x^i} = 0,$$

$$R^2 = \Phi_t + \frac{\partial(u^i \Phi)}{\partial x^i} = 0,$$

$$R^3 = c_t + u^i \frac{\partial c}{\partial x^i} - cn_i n_j \frac{\partial u^i}{\partial x^j} - \alpha(\delta_{ij} - n_i n_j) \frac{\partial}{\partial x^j} \left((\delta_{ik} - n_i n_k) \frac{\partial c}{\partial x^k} \right) = 0.$$

Principal Result 2 (multipliers)

$$\Lambda^1 = \Phi \mathcal{F}(\Phi) |\nabla \Phi|^{-1} \left(\frac{\partial}{\partial x^j} \left(c \frac{\partial \Phi}{\partial x^j} \right) - cn_i n_j \frac{\partial^2 \Phi}{\partial x^i \partial x^j} \right),$$

$$\Lambda^2 = -\mathcal{F}(\Phi) |\nabla \Phi|^{-1} \left(\frac{\partial}{\partial x^j} \left(c \frac{\partial \Phi}{\partial x^j} \right) - cn_i n_j \frac{\partial^2 \Phi}{\partial x^i \partial x^j} \right),$$

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Governing equations ($\alpha = 0$)

$$R^1 = \frac{\partial u^i}{\partial x^i} = 0,$$

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Principal Result 2 (divergence expressions)

- An infinite family of conservation laws:

$$\frac{\partial}{\partial t} (c \mathcal{F}(\Phi) |\nabla \Phi|) + \frac{\partial}{\partial x^i} (A^i \mathcal{F}(\Phi) |\nabla \Phi|) = 0,$$

where

$$A^i = cu^i - \alpha \left((\delta_{ik} - n_i n_k) \frac{\partial c}{\partial x^k} \right), \quad i = 1, 2, 3,$$

and \mathcal{F} is an arbitrary sufficiently smooth function.

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Simplest conservation law with c -dependence

- Can take $\mathcal{F}(\Phi) = 1$:

$$\frac{\partial}{\partial t} (c |\nabla \Phi|) + \frac{\partial}{\partial x^i} (A^i |\nabla \Phi|) = 0.$$

- Surfactant dynamics equations can be written in a **fully conserved form**.

- 1 Fluid Dynamics Equations
- 2 CLs of Constant-Density Euler and N-S Equations
- 3 CLs of Helically Invariant Flows
- 4 CLs of An Inviscid Model in Gas Dynamics
- 5 CLs of a Surfactant Flow Model
- 6 Discussion**

- Fluid & gas dynamics: a large number of general and specific models exist.
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 - single and multi-phase;
 - non-Newtonian;
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- **A lot remains to be discovered!**
- Further applications: numerical simulations; development of specialized numerical methods, etc.



Batchelor, G.K. (2000).

An Introduction to Fluid Dynamics, Cambridge University Press.



Müller, P. (1995).

Ertel's potential vorticity theorem in physical oceanography.
Reviews of Geophysics **33**(1), 67–97.



Kallendorf, C., Cheviakov, A.F., Oberlack, M., and Wang, Y. (2012).

Conservation Laws of Surfactant Transport Equations. *Phys. Fluids* **24**, 102105.



Kelbin, O., Cheviakov, A.F., and Oberlack, M. (2013).

New conservation laws of helically symmetric, plane and rotationally symmetric viscous and inviscid flows. *J. Fluid Mech.* **721**, 340–366.



Cheviakov, A.F., Oberlack, M. (2014).

Generalized Ertels theorem and infinite hierarchies of conserved quantities for three-dimensional time-dependent Euler and NavierStokes equations. *J. Fluid Mech.* **760**, 368–386.